

Lecture 5: Conformal Hilbert Space I

- What is the use of the LEC formalism
 - Simplicity for intuition and numerics
 - Potential for unifying CF theory and CFT / model Hamiltonian approach

Jain state at $2/5$

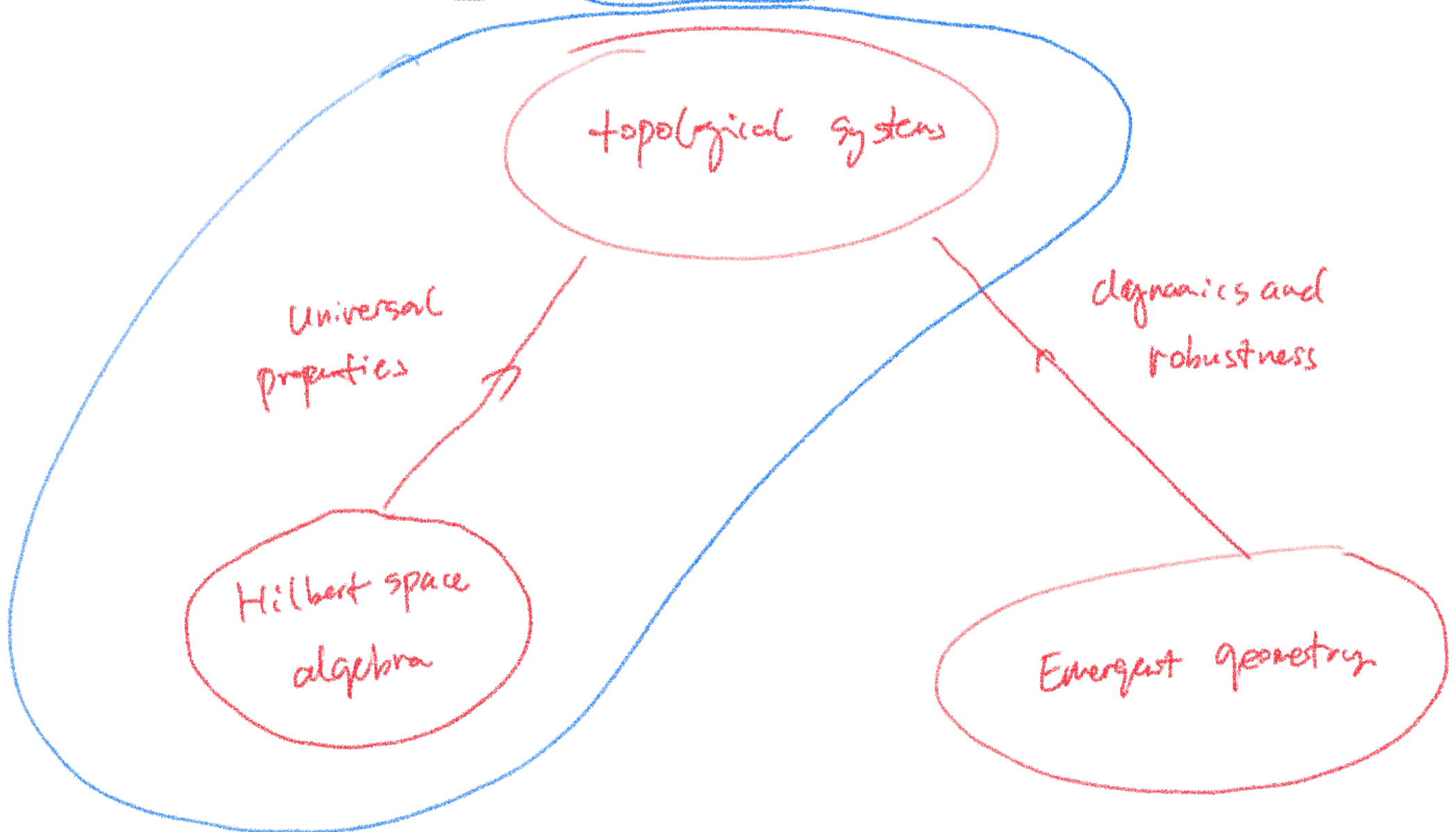
PRB 100, 245303

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arXiv: 2201.00020

- FQHE / IQHE from conformal Hilbert spaces

Topological properties of FQHE fundamentally emerge from Hilbert space truncation



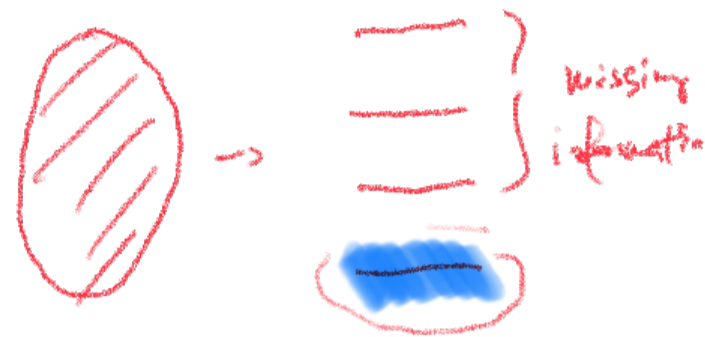
- Full Hilbert space ($B=0$)

... on electrons

- Elementary particles are even -

- Fermions, carrying a charge e , and spin $\frac{1}{2}$
- Effective mass m^* or Fermi velocity V_F^*
- point particles ($r < 10^{-22}$ m)

• A single Landau level (LL) $B \rightarrow \infty$



- A sub-Hilbert space

- Elementary particles are "quasiparticles"

- Fermions, carrying a charge of e , and spin $\frac{1}{2}$

- No dispersion, $m^* \rightarrow \infty$ or $V_F^* \rightarrow 0$

- NOT point particles (occupy a finite area $2\pi l_B^2$)

$$[\vec{R}^x, \vec{R}^y] = -i l_B^2$$

↓
a loss of information
due to LL projection

- A Hilbert space with conformal symmetry

↳ angle-preserving

- $P_\mu = -i \partial_\mu$ (translation)

$D = -i x^\mu \partial_\mu$ (dilatation)

$L_{\mu\nu} = x_\mu p_\nu - x_\nu p_\mu$ ($\vec{r} \times \vec{p}$) (rotation)

$K_\mu = -i (2x_\mu x^\nu \partial_\nu - x^2 \partial_\mu)$ (SCT)

$$L^x = L_{yz}$$

$$L^y = L_{zx}$$

$$L^z = L_{xy}$$

there is a hidden metric

If we define $T_{\mu\nu} = L_{\mu\nu}$ $\mu, \nu = 1, 2, 3$
 x, y, z

$T_{-1,0} = D$

$$\bar{J}_{0,\mu} = \frac{1}{2} (P_\mu + K_\mu)$$

$$J_{-1,\mu} = \frac{1}{2} (P_\mu - K_\mu)$$

$$J_{ab} = -J_{ba}$$

$$\Rightarrow [J_{ab}, J_{cd}] = i [\eta_{ad} J_{bc} + \eta_{bc} J_{ad} - \eta_{ac} J_{bd} - \eta_{bd} J_{ac}]$$

$$\eta_{ab} = \text{diag}(-1, 1, \dots, 1)$$

Conformal group in d -dimension

$$\cong \underline{\underline{SO(d+1, 1)}} \quad (-1, 1, \dots, 1) \\ = \underbrace{\hspace{2cm}}_{d+1}$$

• for $d=2$, $SO(3, 1)$ is the Lorentz group

• Conformal Symmetry as a "glorified" version of $SU(2)$ algebra (rotation)

- In two-dimension, we are interested in local conformal symmetry

• The Virasoro algebra

$$[\hat{L}_n, \hat{L}_m] = (n-m) \hat{L}_{m+n} + \frac{c}{12} n(n^2-1) \delta_{m+n,0}$$

n, m are integers

→ central charge

↓
a quantum anomaly

- quantization from macroscopic scale

"softly" break conformal symmetry

★ The composition of conformal Hilbert space

- A highest density state (Fock ground state)



$$\begin{aligned} | \psi_{h_0} \rangle &\rightarrow \hat{L}_0 | \psi_{h_0} \rangle = h_0 | \psi_{h_0} \rangle \\ \underline{\underline{=}} \quad \hat{L}_{n>0} | \psi_{h_0} \rangle &= 0 \quad (\text{HW condition}) \end{aligned}$$

present for
non-Abelian
FQH

- A collection of highest weight states
 - $\{ | \psi_{h_i} \rangle \} \quad i=1,2,\dots \dots n_h \quad h_i > h_0$
 - $\underline{\underline{\{ | \psi_{h_i} \rangle \}}}$
 - $\rightarrow \hat{L}_0 | \psi_{h_i} \rangle = h_i | \psi_{h_i} \rangle$
 - $\hat{L}_{n>0} | \psi_{h_i} \rangle = 0$
 - (FQH quasiparticle states)

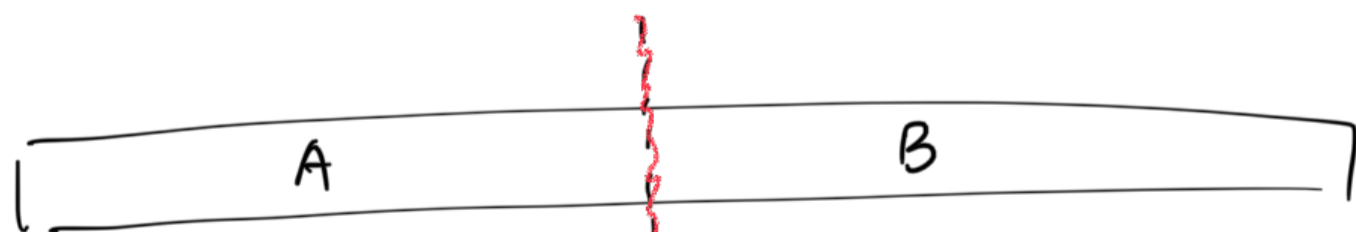
- All other "secondary states" (quasihole states)

- Generated by applying

$$\hat{L}_{n<0} \text{ onto } \underline{\underline{\{ | \psi_{h_i} \rangle \}}}, \underline{\underline{\{ | \psi_{h_0} \rangle \}}}$$

- The bulk-edge correspondence

Ground state with Ne



\rightarrow a virtual edge = a physical edge

$$\rho_A = \sum_i \lambda_i | \psi_{iA} \rangle \langle \psi_{iA} |$$

\downarrow
this basis

spans the same conformal Hilbert space at $N_e/2$ electrons in the limit of $N_e \rightarrow \infty$

- The Hilbert space of a single LL is a Conformal Hilbert space

- It is the null space of the kinetic Hamiltonian

- It is the space spanned by states satisfying
 $\hat{C} = [1, 1, 1]$

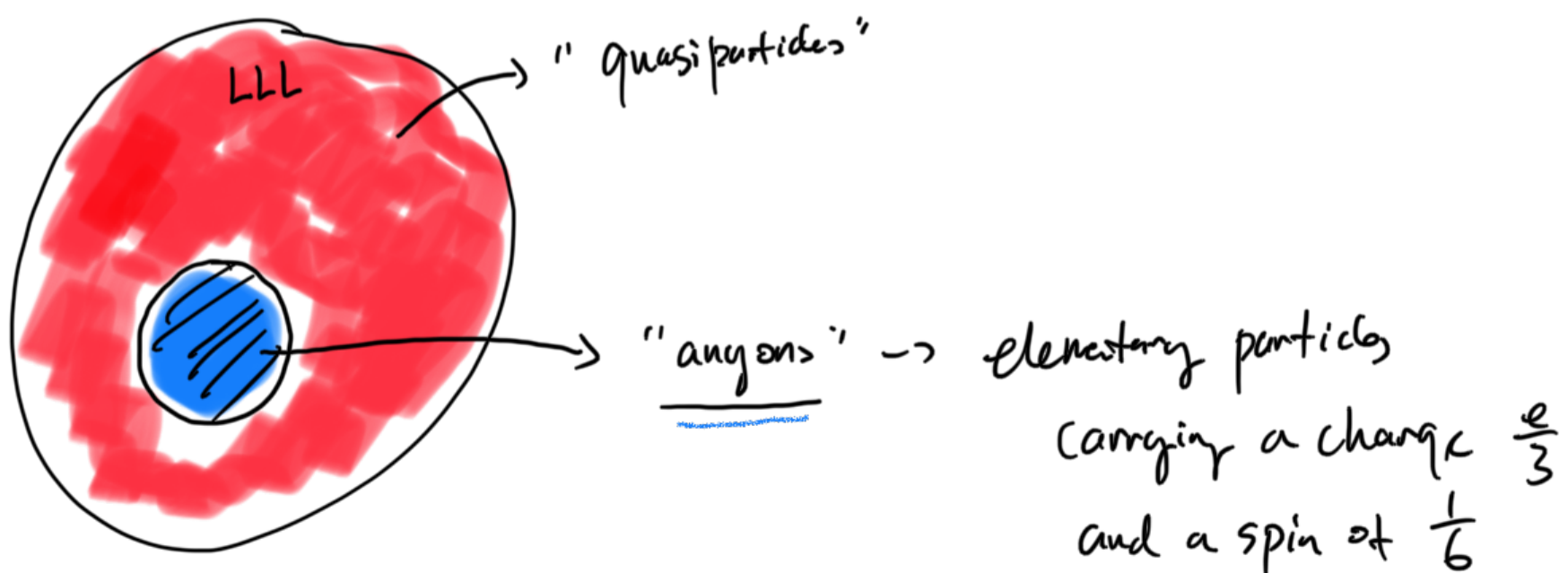
- Other examples of Conformal Hilbert spaces

- Null space of model Hamiltonian ✓

- Space spanned by a particular family of Jack polynomials satisfying certain admissible rules ✓

- Space spanned by states satisfying LEC

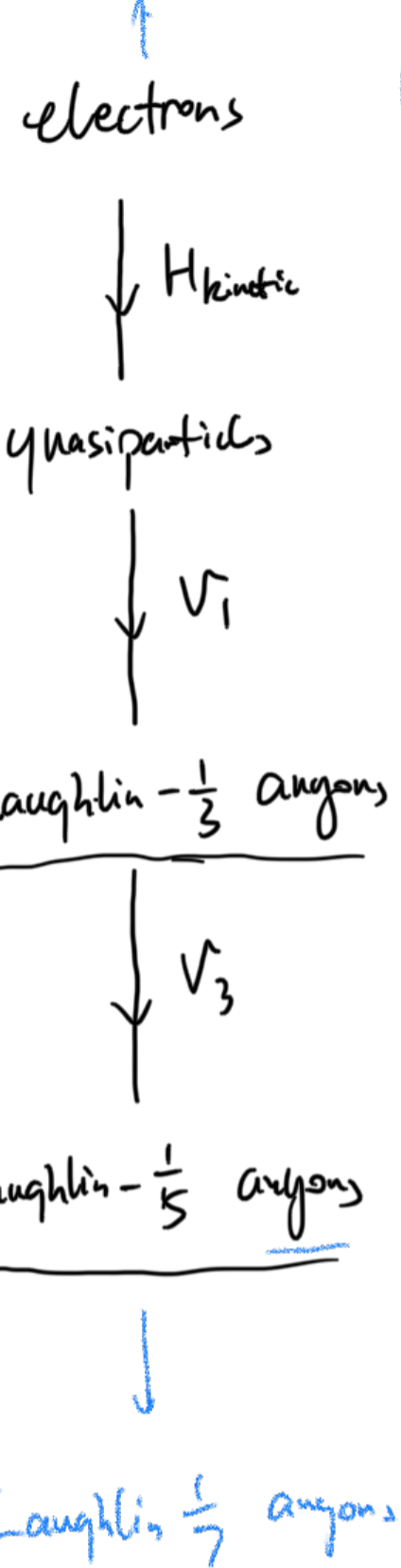
- Hierarchy of CHS:



• → null space of V_i

quarks

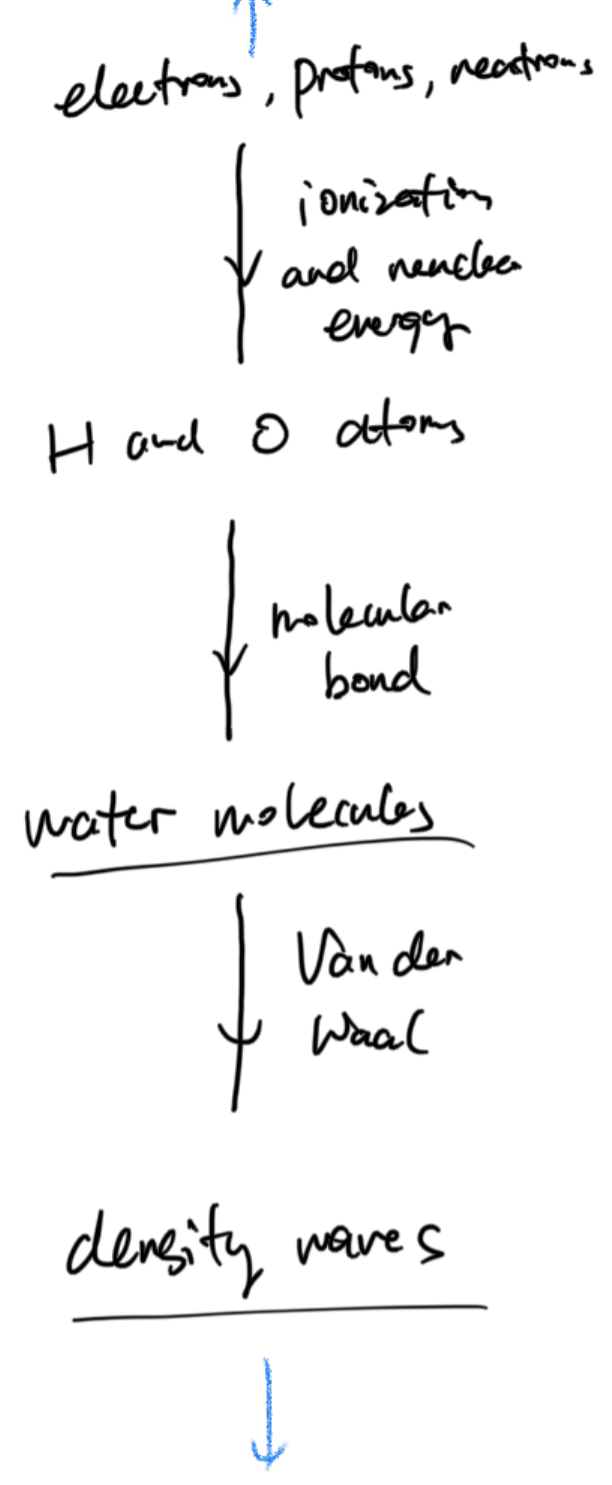
↑ E



$H = V_1$
 $\hat{C} = [2, 1, 2]$

$H = V_1 + V_3$
 $\hat{C} = [3, 1, 3]$

≈



- Let $\mathcal{H}_{\frac{1}{2m+1}}$ be the null space of Laughlin $-\frac{1}{2m+1}$

$\mathcal{H}_{\frac{1}{2m+1}} \subset \mathcal{H}_{\frac{1}{2m'+1}}$ for $m > m'$

- $\mathcal{H}_{\frac{1}{3}} \subset \mathcal{H}_{\text{Moore-Read}}$
 $\hat{C} = [2, 1, 2]$ $\hat{C} = [3, 2, 3]$

- The Read-Rezayi Series

$\mathcal{H}_{k+1} \supset \mathcal{H}_k$

↓

$\hat{C} = [k+1, k, k+1]$

$\mathcal{H}_{k=1} \rightarrow \text{Laughlin } \frac{1}{3}$

$\mathcal{H}_{k=2} \rightarrow \text{Moore-Read } \frac{2}{5}$

$\mathcal{H}_{k=3} \rightarrow \text{Fibonacci } \frac{3}{7}$

- The Conformal Hilbert space formalism is a unifying picture for both ZQH / FQH

- Each Conformal Hilbert space has a vacuum \Rightarrow the ground state of FQHE

- Insertion of magnetic fluxes creates elementary particles \rightarrow anyons

- All physics of FQH / ZQH are determined by the algebraic properties of CHS.

- Such physics is realised when the Hamiltonian creates an energy gap $\gg k_B T$ or disorder between $|\psi\rangle \in \mathcal{H}_{\text{CHS}}$ and $|\psi\rangle \notin \mathcal{H}_{\text{CHS}}$
