

Universal entropy of conformal critical theories on a Klein bottle: a quantum Monte Carlo study

Wei Tang (唐维)

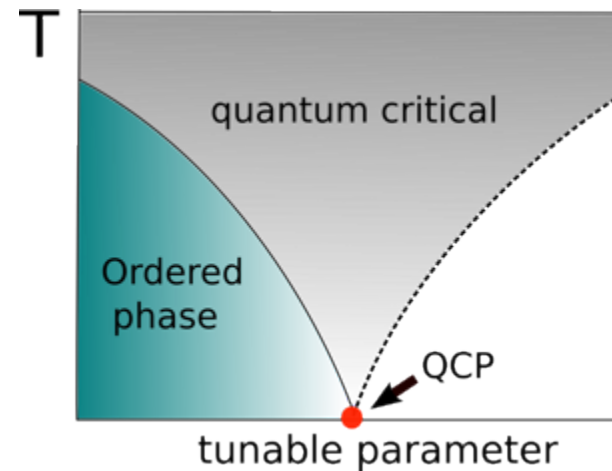
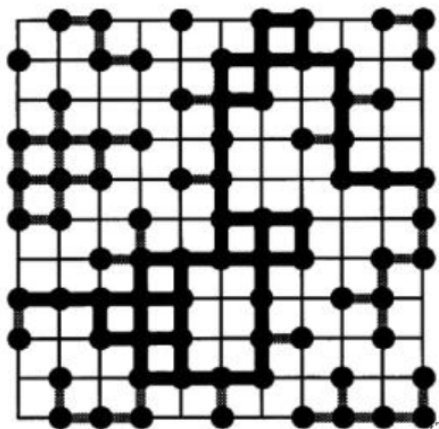
ICQM PKU

Collaborators

- Lei Wang, IOP CAS
- Hong-Hao Tu, TU Dresden
- Lei Chen, Wei Li, Beihang University
- X.-C. Xie, ICQM PKU

Critical phenomena and conformal field theory

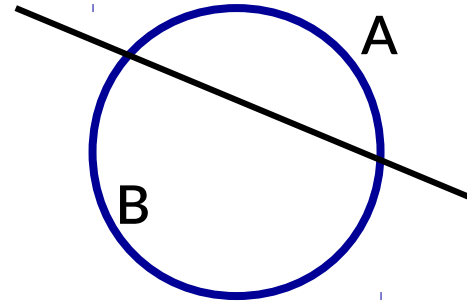
- Conformal invariance at the critical points
- In 2D, the conformal invariance gives very strong constraints
→ critical points of 2D classical models or 1D quantum models



Address the underlying CFT of a microscopic system

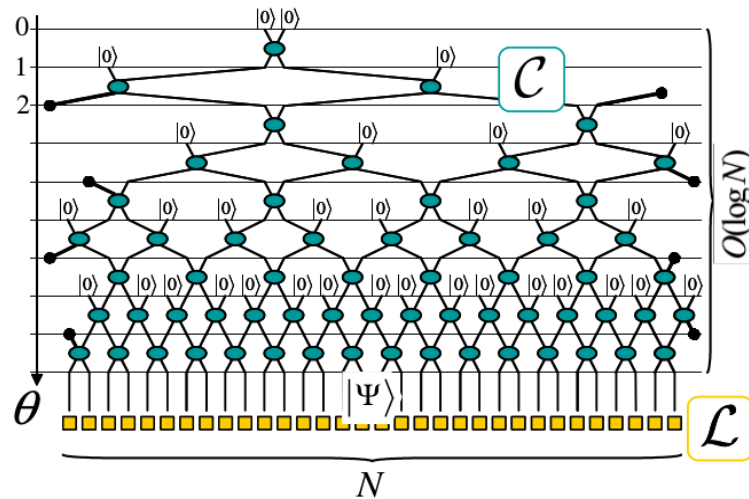
- Entanglement entropy

$$S_A \sim \frac{c}{3} \ln L$$



Calabrese and Cardy,
J. Stat. Mech. 2004,
J. Phys. A 2009

- MERA



Vidal, PRL 2008

This work → Klein bottle entropy: a universal quantity as a characterization of the CFT

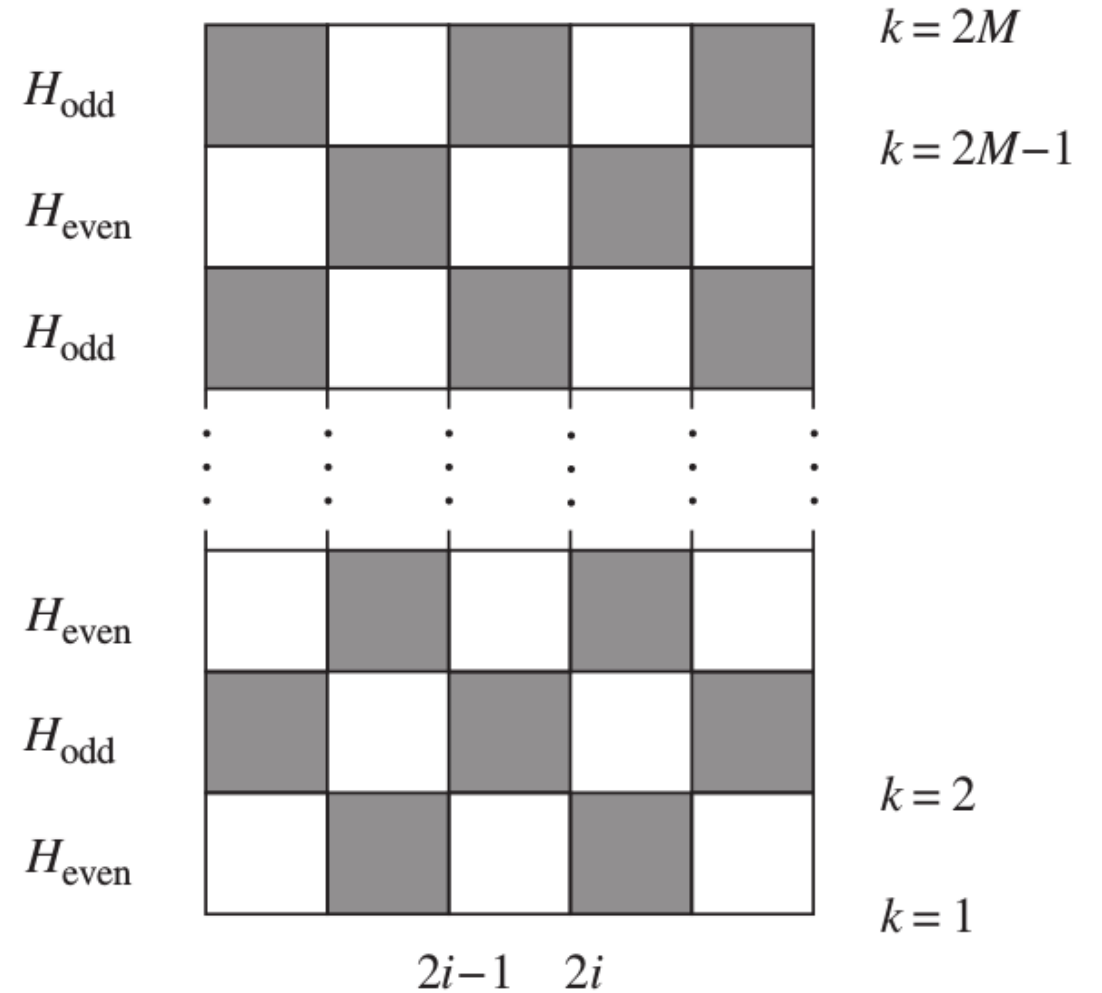
Outline

- Critical phenomena and conformal field theory
- **Universal boundary entropies on various manifolds**
- Calculations in lattice models
- Models and results
- Summary

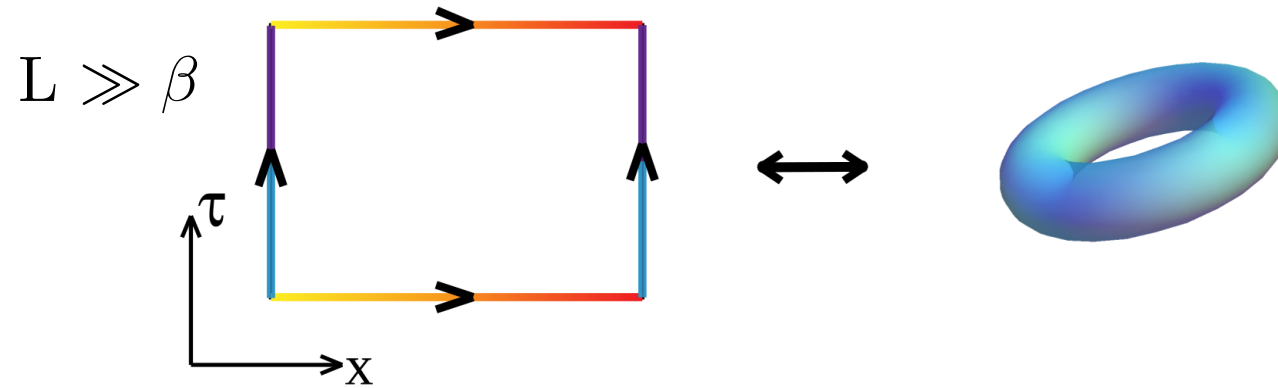
The 1D quantum system at finite temperature

$$\begin{aligned}
 Z^{\mathcal{T}} &= \text{tr}(e^{-\beta H}) \\
 &= \text{tr}(e^{-\Delta\tau H_{\text{even}}} e^{-\Delta\tau H_{\text{odd}}} \dots e^{-\Delta\tau H_{\text{odd}}})
 \end{aligned}$$

The partition function of a 1D quantum system with periodic boundary conditions lives on a torus



Torus partition function (of a long strip)



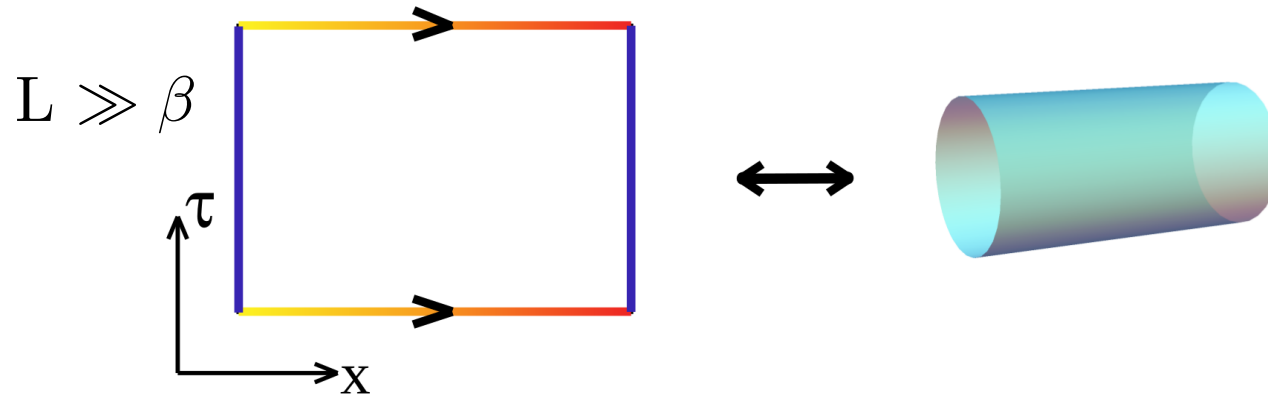
$$\ln Z^{\mathcal{T}} = -f_0 \beta L + \frac{\pi c}{6\beta v} L + O\left(\frac{1}{\beta^2}\right)$$

- f_0 the bulk free energy density
- $\frac{\pi c}{6\beta v} L$ the bulk entropy (Prediction of CFT)

Bloete, Cardy and Nightingale, PRL 1986

Affleck, PRL 1986

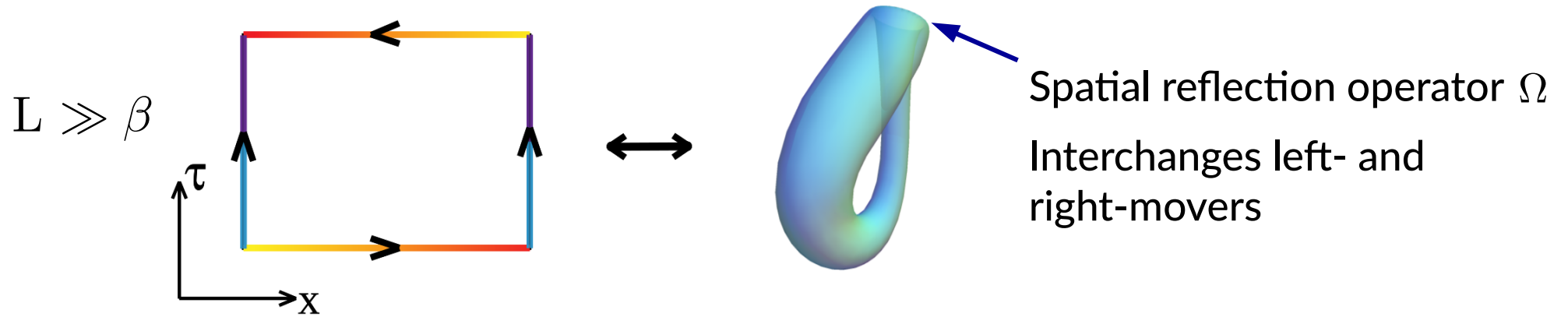
Cylinder partition function (of a long strip)



$$\ln Z^c = -f_0 \beta L + \frac{\pi c}{6\beta v} L + S_{\text{AL}} - f_b \beta + O\left(\frac{1}{\beta^2}\right)$$

- f_b the surface free energy density
- S_{AL} the universal Affleck-Ludwig entropy (Prediction of CFT)

Klein bottle partition function (of a long strip)



$$\ln Z^{\mathcal{K}} = -f_0 \beta L + \frac{\pi c}{24 \beta v} L + S_{\text{KB}} + O\left(\frac{1}{\beta^2}\right)$$

- Different coefficient of the bulk entropy
- An additional universal term S_{KB} resembling the AL entropy
→ **The Klein bottle entropy**

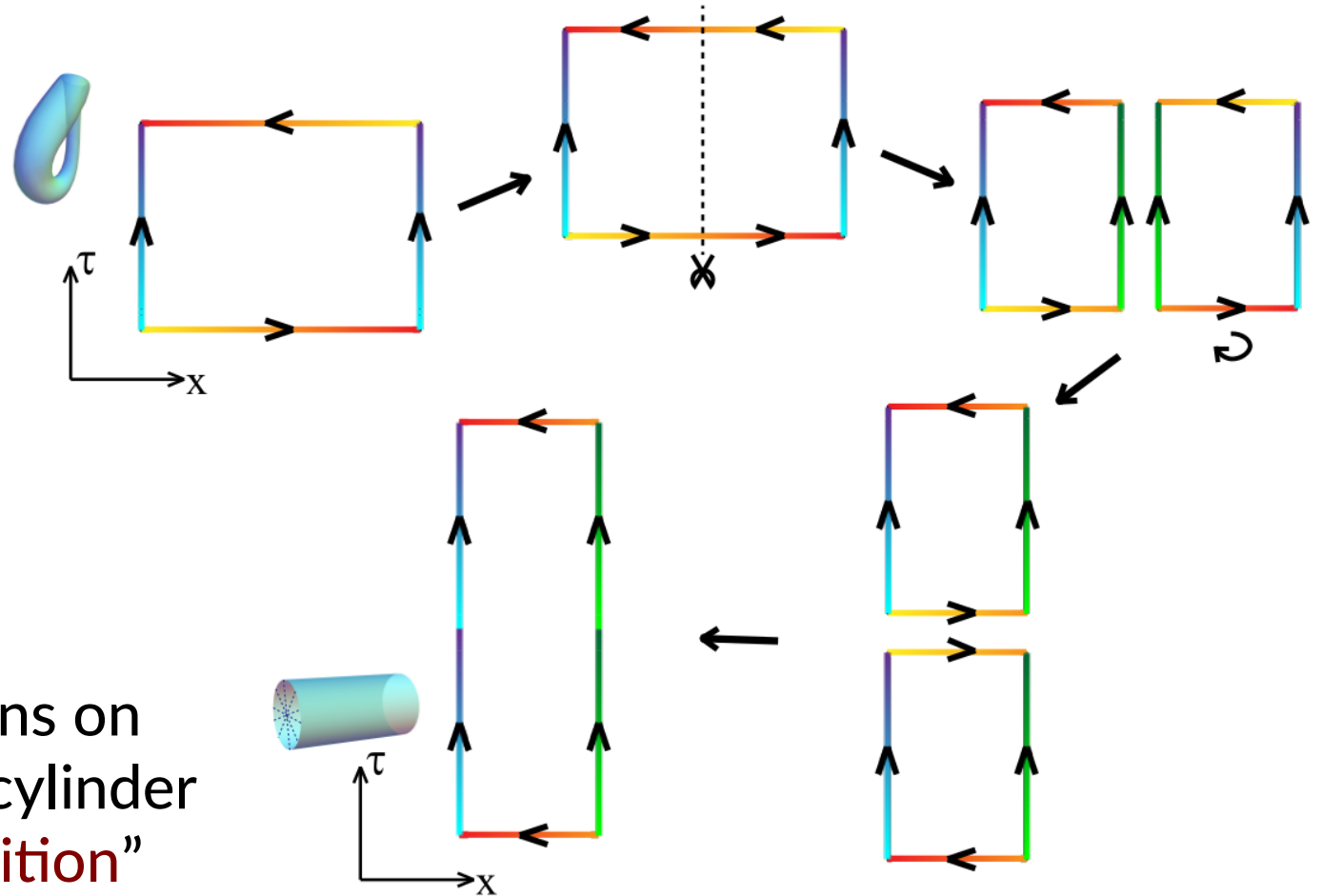
Transformation of the Klein bottle

Klein bottle (L, β)

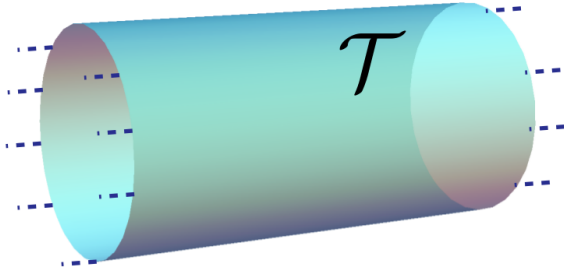
$$\downarrow \frac{\pi c}{24\beta v} L = \frac{\pi c}{6(2\beta)v} \frac{L}{2}$$

cylinder $(L/2, 2\beta)$

long-range interactions on the boundary of the cylinder
 \rightarrow “a boundary condition”

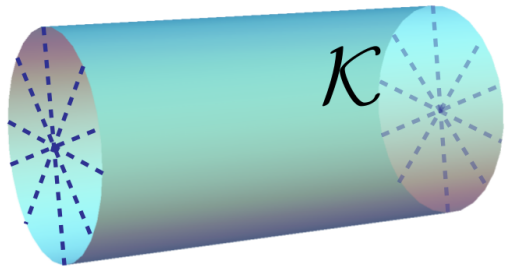


Klein bottle entropy: a boundary entropy



$$\ln Z^{\mathcal{T}}(L, \beta) = -f_0\beta L + \frac{\pi c}{6\beta v}L + O\left(\frac{1}{\beta^2}\right)$$

emerges from the new
boundary condition



$$\ln Z^{\mathcal{K}}(2L, \beta/2) = -f_0\beta L + \frac{\pi c}{6\beta v}L + \boxed{S_{\text{KB}}} + O\left(\frac{1}{\beta^2}\right)$$

➔ $S_{\text{KB}} = \ln \left[\frac{Z^{\mathcal{K}}(2L, \beta/2)}{Z^{\mathcal{T}}(L, \beta)} \right]$

Extracts the ground-state property
from thermal systems

CFT predictions of the Klein bottle entropy

- In rational CFT,

$$S_{\text{KB}} = \ln\left(\sum_a M_{aa} d_a / \mathcal{D}\right)$$

d_a the quantum dimension
of the primary field

Tu, PRL 2017

- In the compactified boson CFT,

$$S_{\text{KB}} = \ln R$$

R the compactification radius
 $\phi = \phi + 2m\pi R, m = 0, \pm 1, \dots$

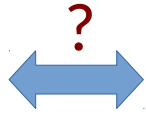
WT, Xie, Wang and Tu, 1805.01300

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Bond-centered lattice reflection

Lattice reflection



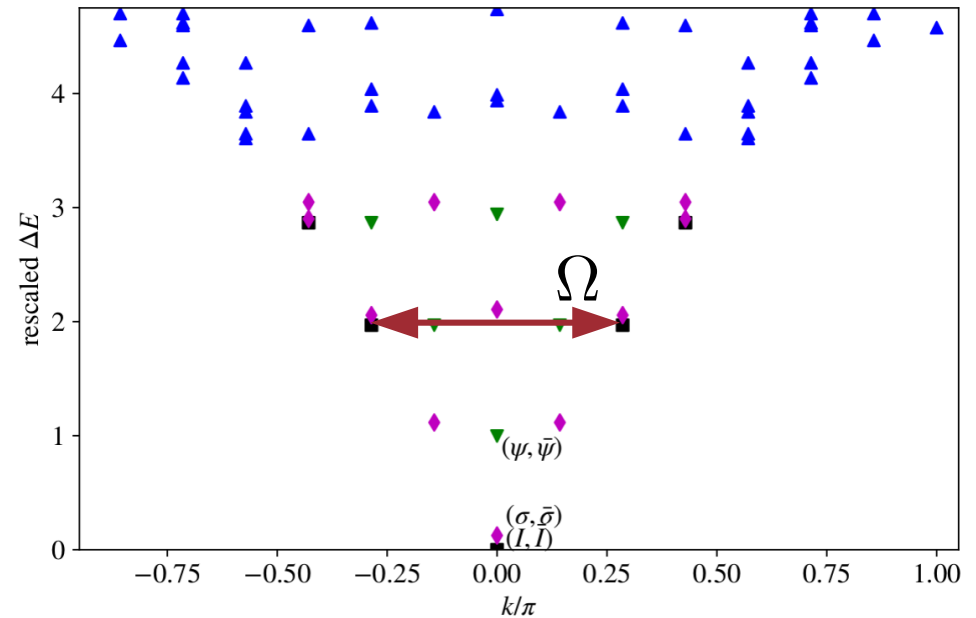
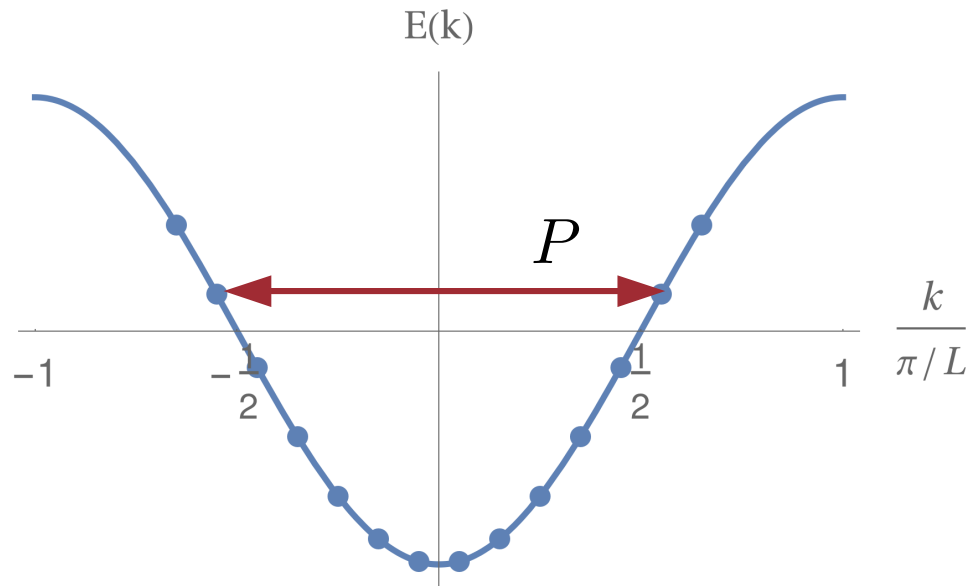
$$P|s_1, \dots, s_L\rangle = |s_L, \dots, s_1\rangle$$

$$Z^K = (Pe^{-\beta H})$$

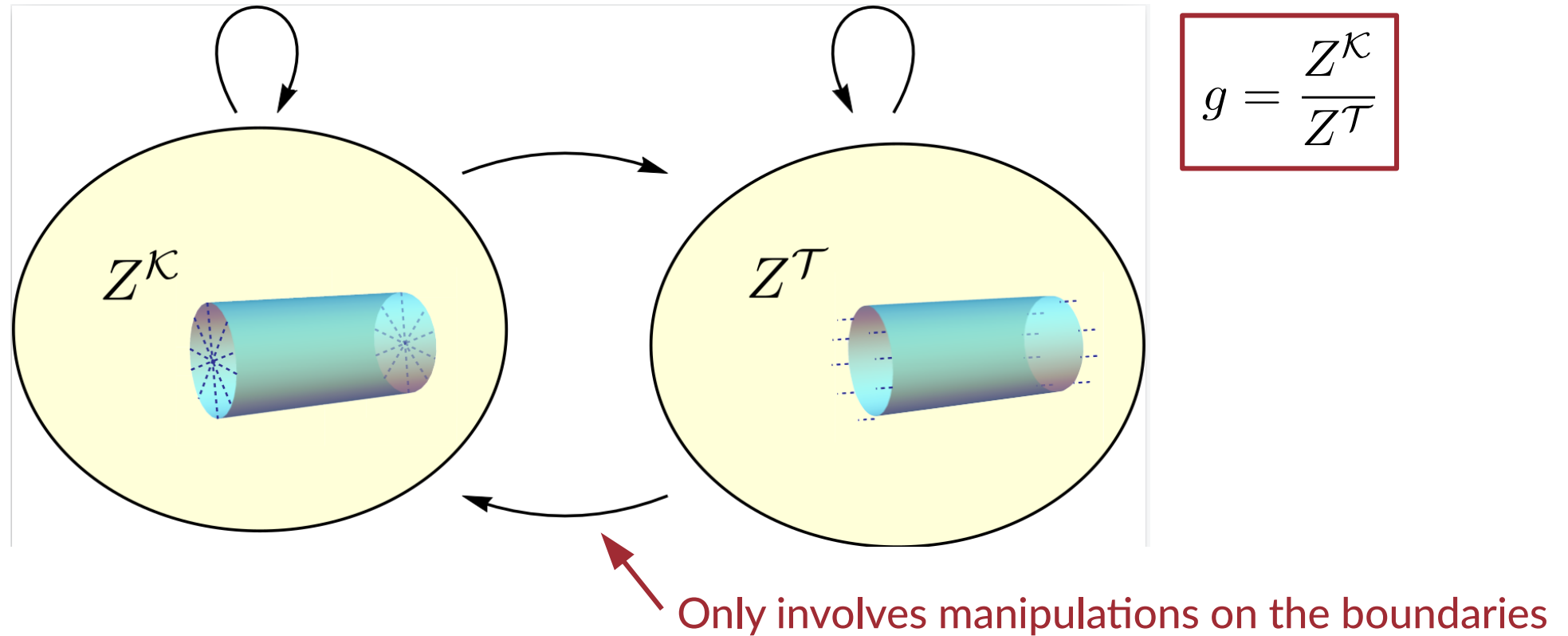
The reflection Ω in the CFT Hilbert space

$$\Omega|\alpha, \bar{\gamma}\rangle = |\gamma, \bar{\alpha}\rangle$$

$$Z^K = (\Omega e^{-\beta H})$$



Extended-ensemble Monte Carlo for partition function ratios



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Spin-1/2 XY chain

$$H = - \sum_{i=1}^L (S_i^x S_{i+1}^x + S_i^y S_{i+1}^y)$$

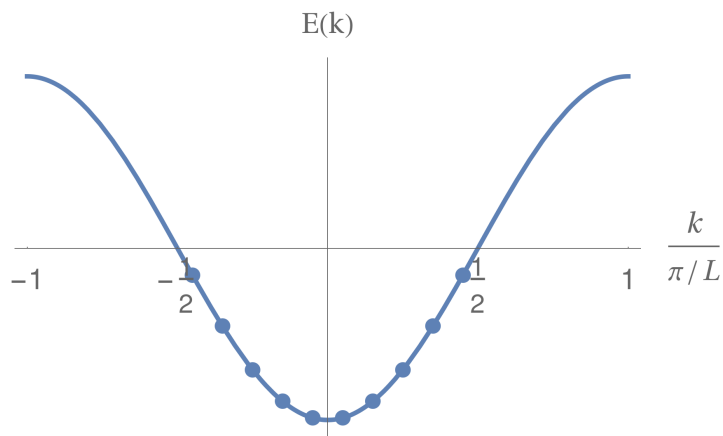
↓ Jordan-Wigner

$$H_{\pm} = -\frac{1}{2} \sum_{i=1}^L (f_i^{\dagger} f_{i+1} + \text{h.c.})$$

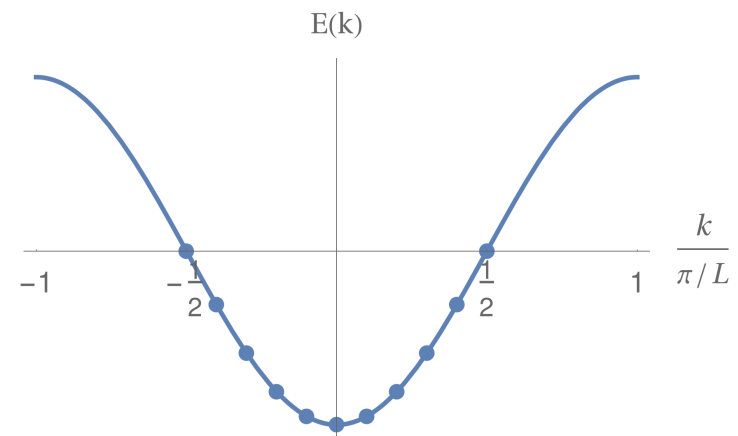
$$= - \sum_k \cos(k) f_k^{\dagger} f_k$$

NS sector: anti-PBC $f_1 = -f_{L+1}$ even number of fermions

R sector: PBC $f_1 = f_{L+1}$ even number of fermions

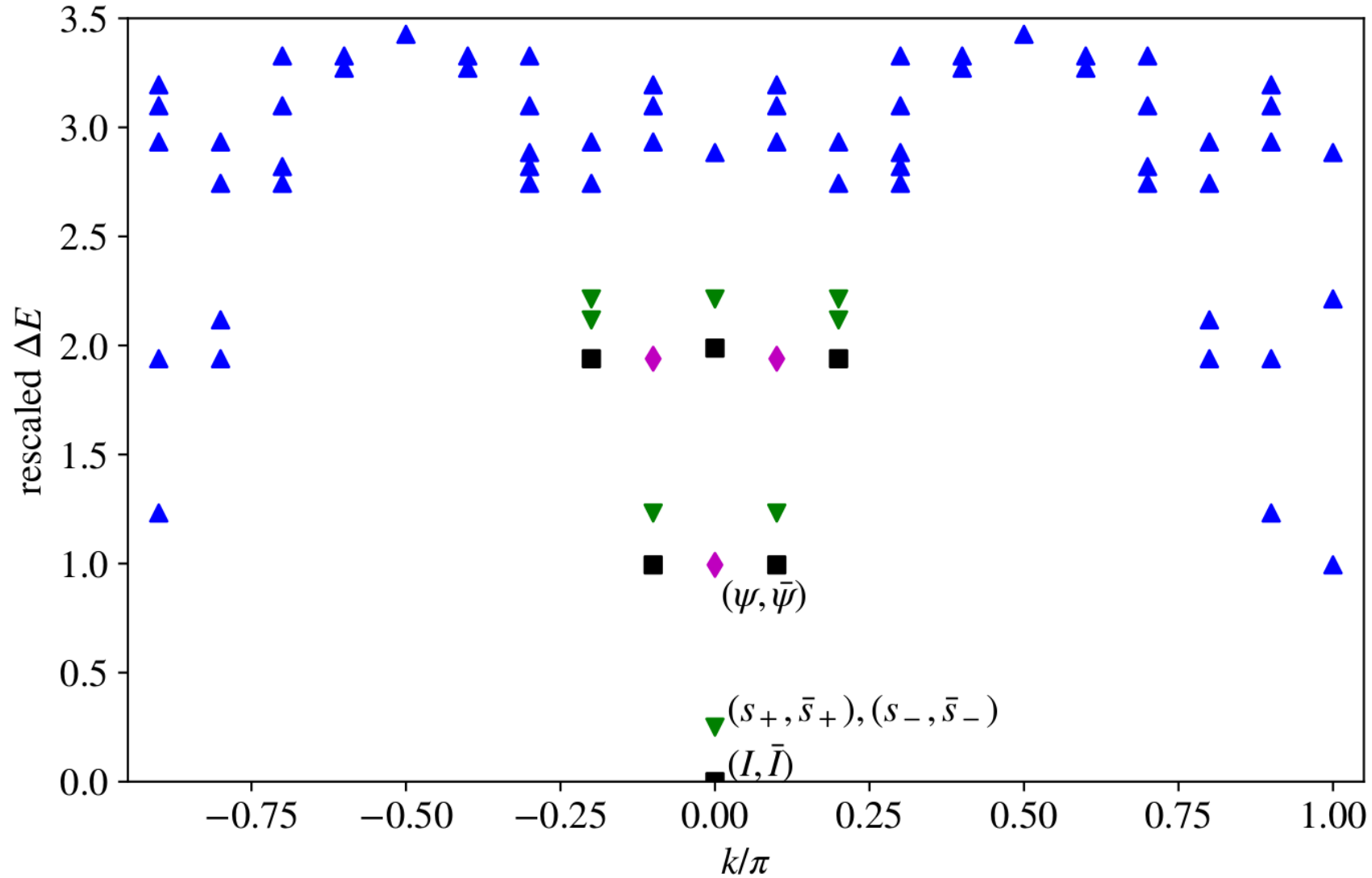


NS sector

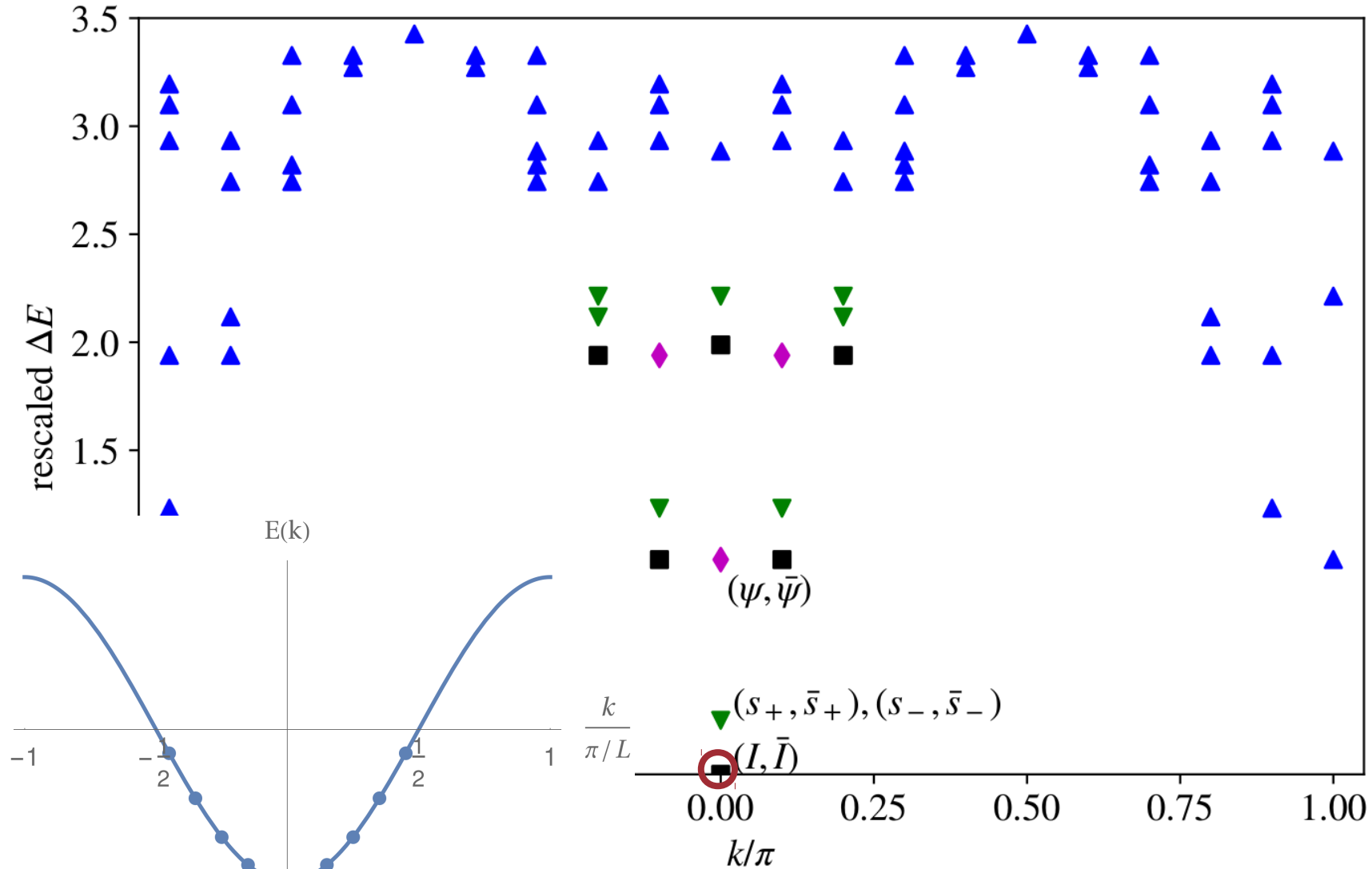


R sector

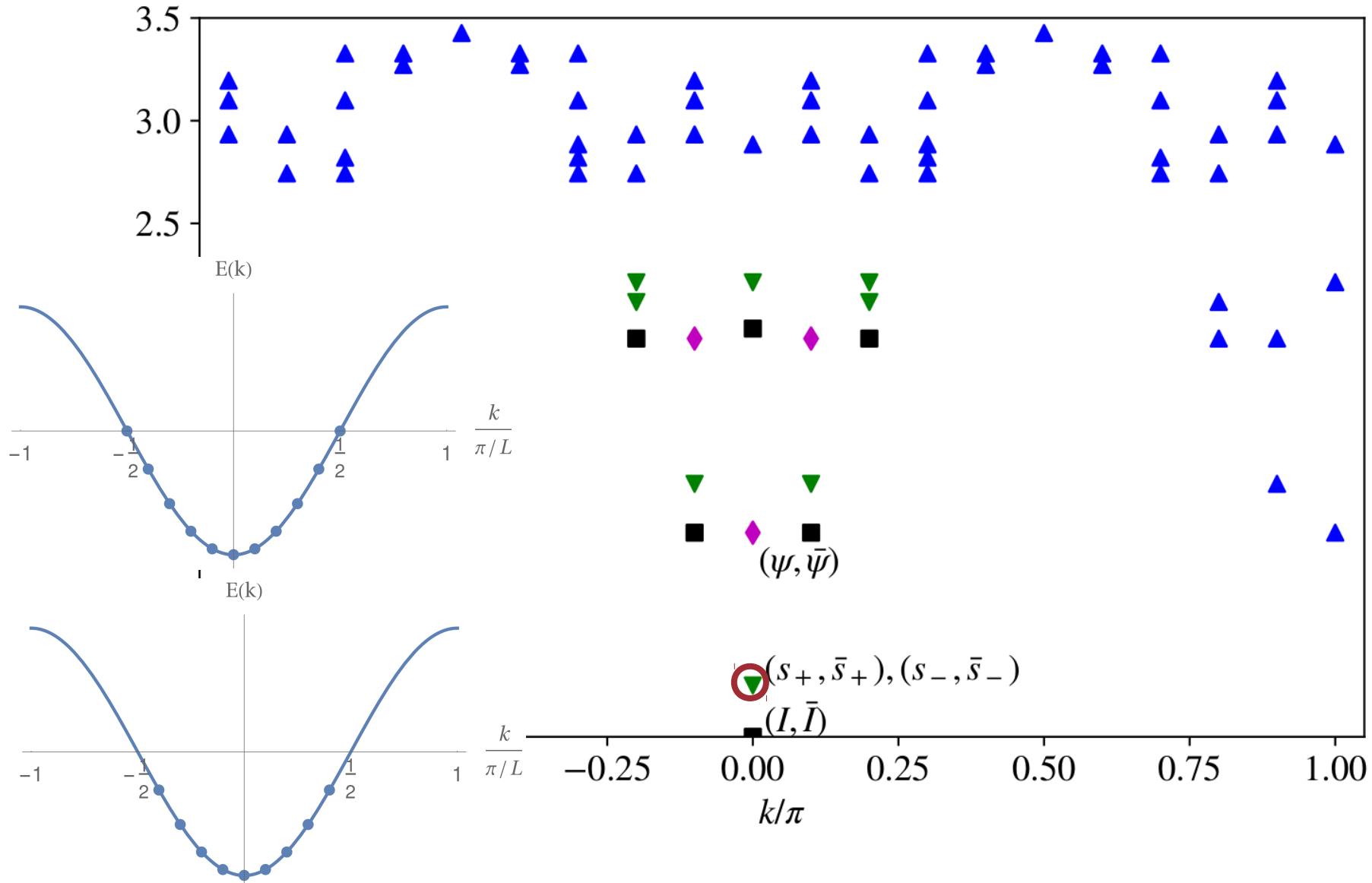
Spin-1/2 XY chain (rational CFT analysis)



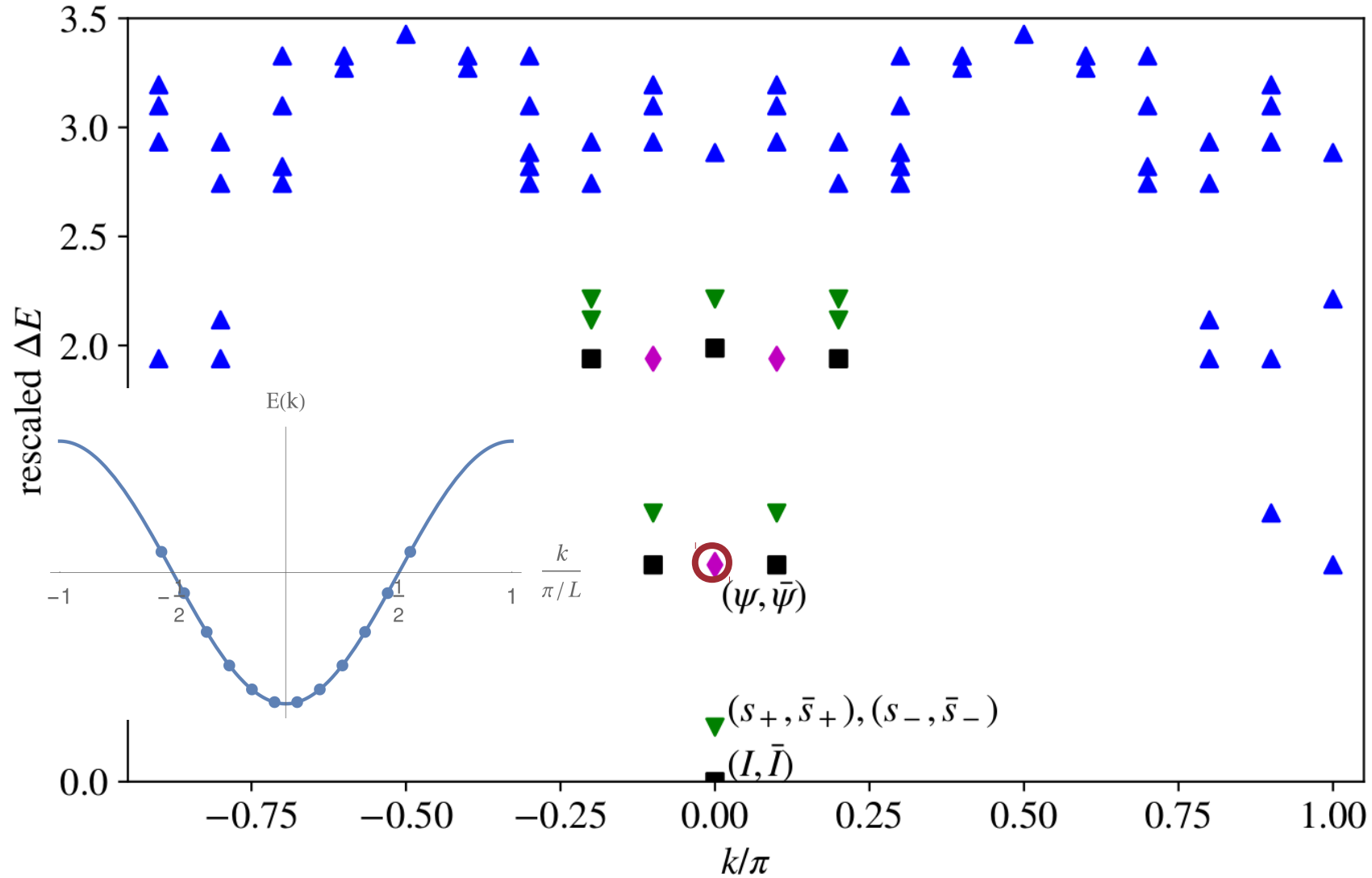
Spin-1/2 XY chain (rational CFT analysis)



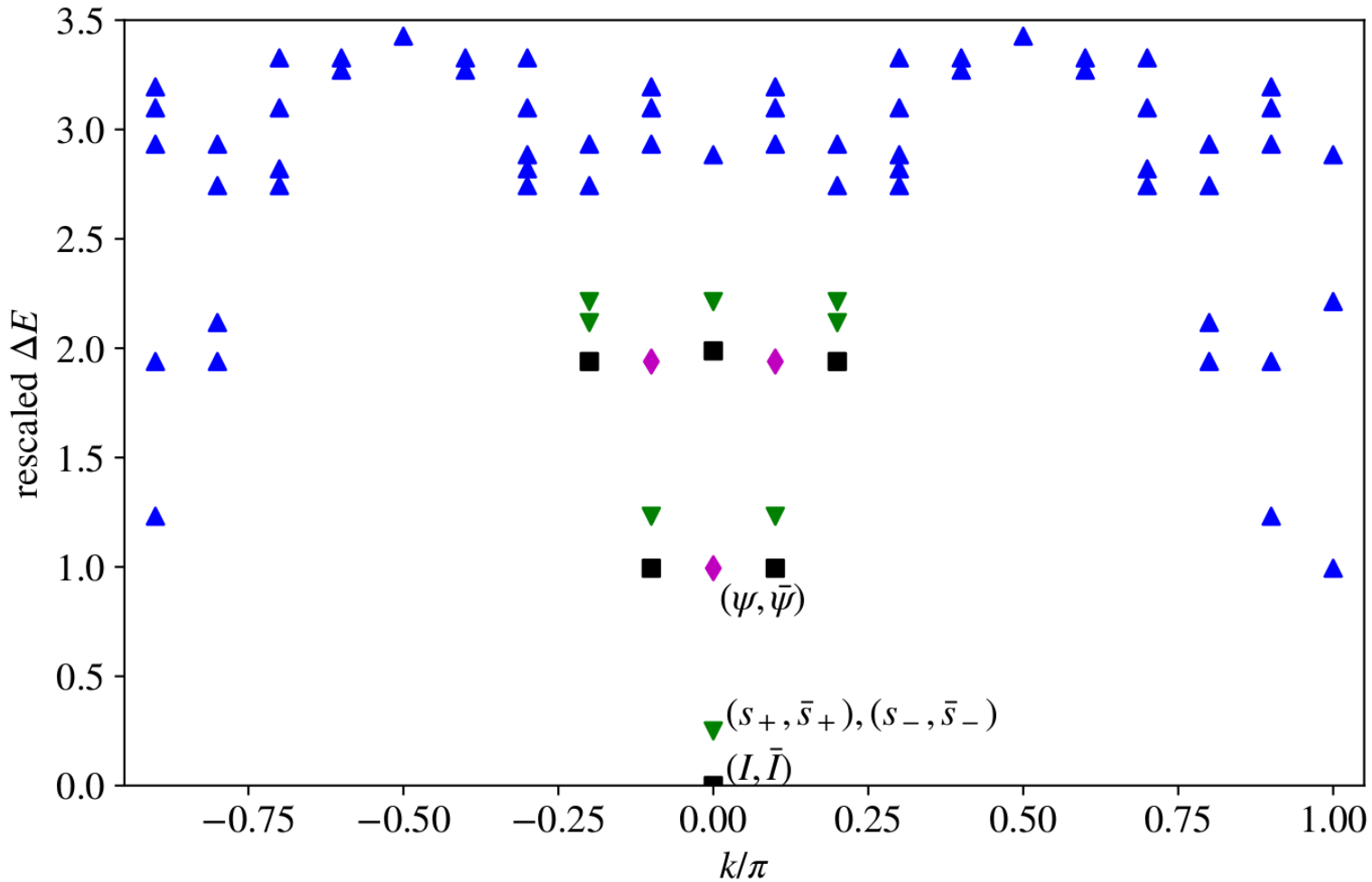
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Spin-1/2 XY chain (rational CFT analysis)



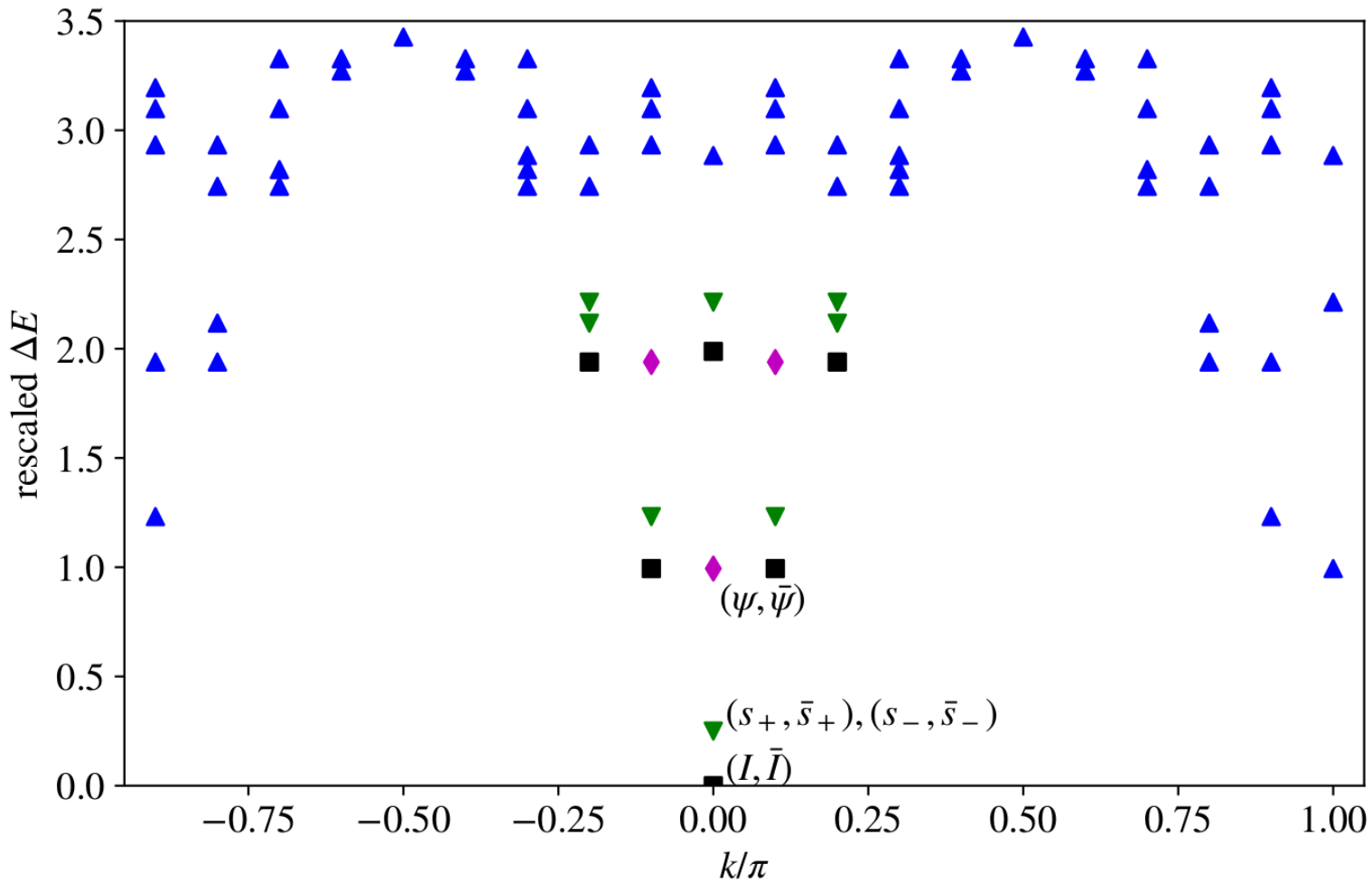
$$d_I = d_{s_+} = d_{s_-} = d_\psi = 1$$

$$\mathcal{D} = \sqrt{1 + 1 + 1 + 1} = 2$$



$$S_{\text{KB}} = \ln \frac{1 + 1 + 1 + 1}{2} = \ln 2$$

Spin-1/2 XY chain (rational CFT analysis)



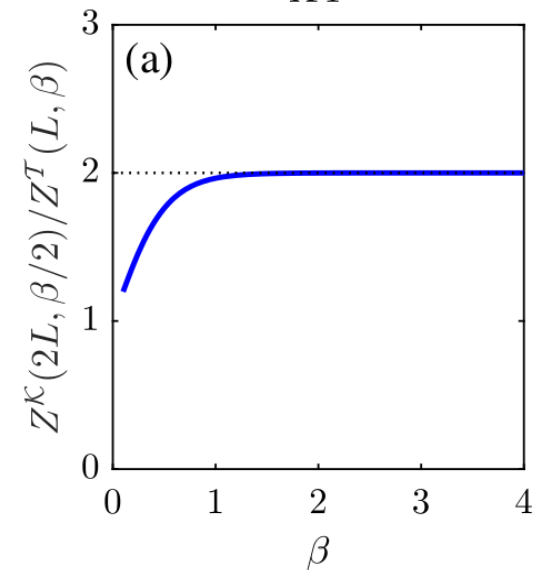
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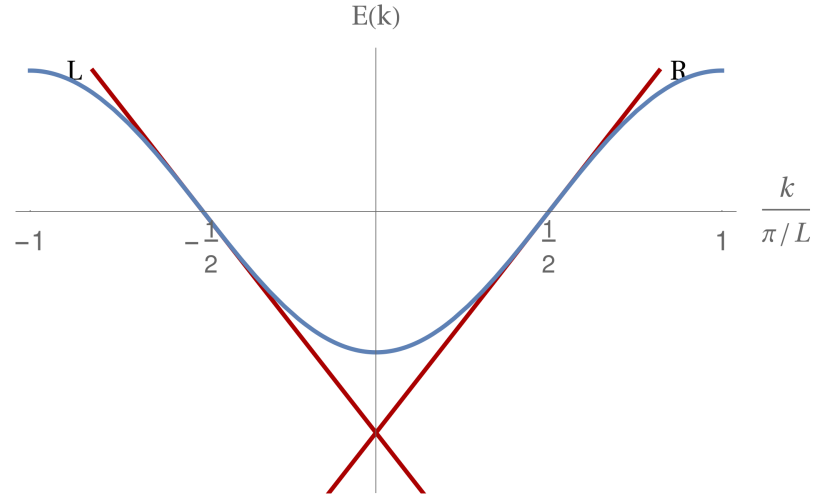
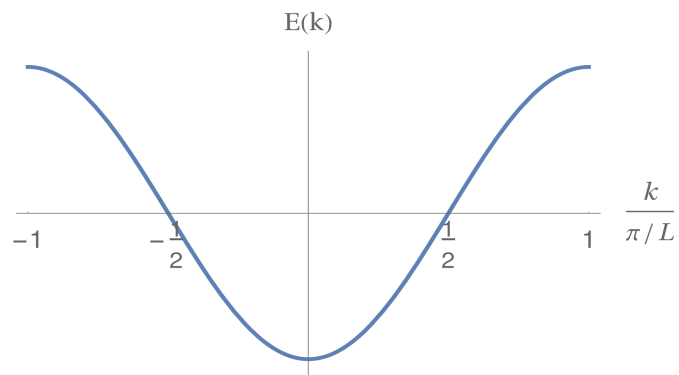


$$S_{\text{KB}} = \ln \frac{1 + 1 + 1 + 1}{2} = \ln 2$$

XY

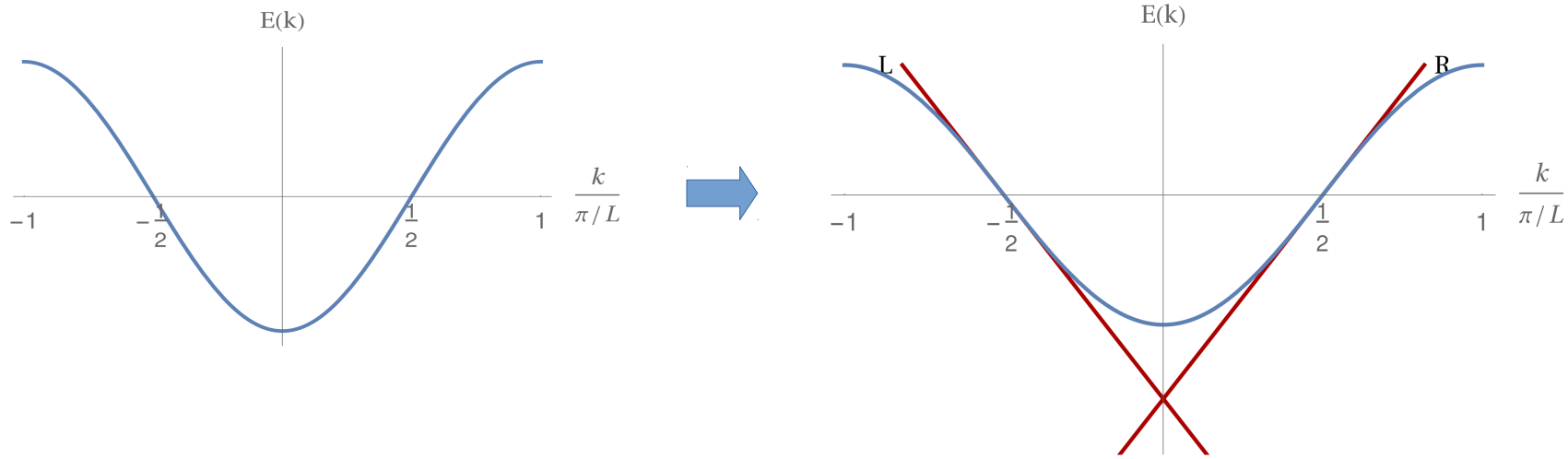


Luttinger liquid description of XY chain



Luttinger Liquid
(spinless, non-interacting)

Luttinger liquid description of XY chain



Luttinger Liquid
(spinless, non-interacting)

The structure of the Hilbert space of LL

Fermi-sea states

$$|n_L, n_R\rangle$$

+

Particle-hole excitations

$$\rho_{L/R, -q} = \frac{2\pi}{L} \sum_k : f_{k+q, L/R}^\dagger f_{k, L/R} :$$

Lattice reflection



$$n_L \leftrightarrow n_R$$

$$\rho_{L/R, -q} \leftrightarrow \rho_{R/L, -q} \text{ (Up to a phase factor)}$$

Spin-1/2 XY chain (compactified boson CFT analysis)

(spinless) Luttinger liquid

↓ bosonization

Compactified free boson

$$\phi(x) = \phi(x + L) + 2\pi m R, m \in \mathbb{Z}$$

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Bosonization identity $\psi_{L/R}(x) = a^{-1/2} e^{\pm i \frac{\pi}{L} x} F_{L/R} e^{-i \phi_{L/R}(x)}$

Left- and right-moving boson:

$$\phi_{L/R}(x) = \sum_{q \neq 0} \frac{i}{q} e^{-a|q|/2} e^{\mp i q x} \rho_{q,L/R} \pm \rho_{0,L/R} x$$

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Responsible for the discontinuity
at the boundary

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Responsible for the discontinuity at the boundary

Dual fields

$$\phi = \phi_L + \phi_R$$

$$\theta = \phi_L - \phi_R$$

“T-duality”

Spin-1/2 XY chain (compactified boson CFT analysis)

(spinless) Luttinger liquid

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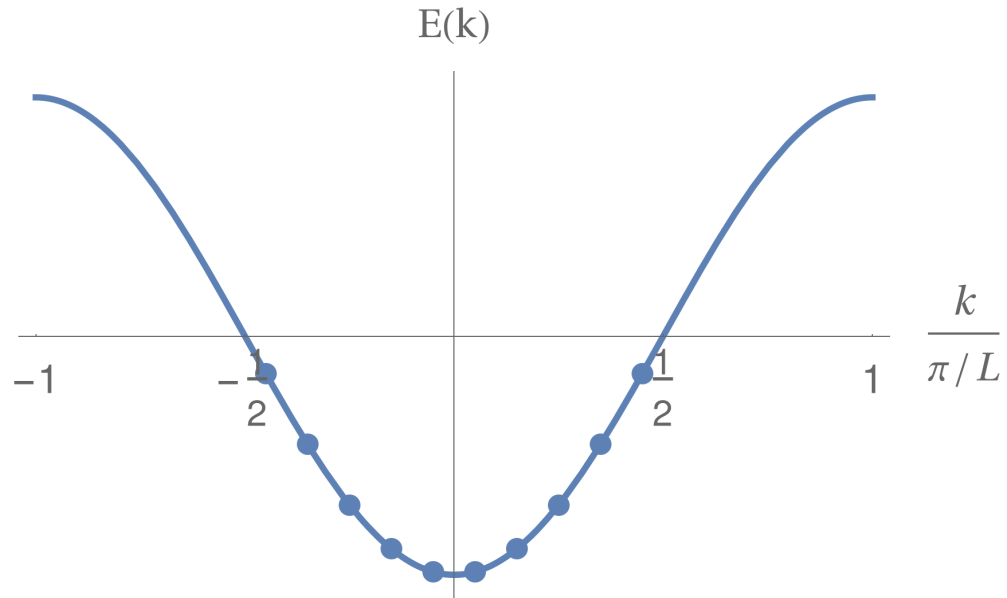
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Responsible for the discontinuity at the boundary

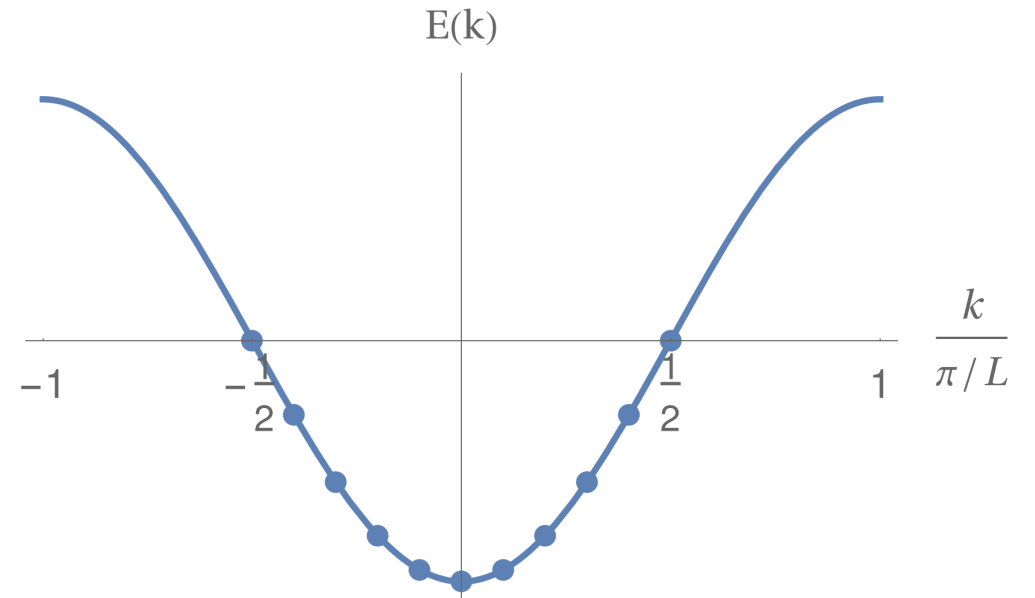
Dual fields

$$\begin{array}{l} \phi = \phi_L + \phi_R \\ \theta = \phi_L - \phi_R \end{array} \quad \begin{array}{l} \nearrow \\ \rightarrow \end{array} \quad \begin{array}{l} \rho_{0,L/R} = \frac{2\pi}{L} n_{L/R} \\ \phi(L) - \phi(0) = 2\pi(n_L - n_R) \\ \theta(L) - \theta(0) = 2\pi(n_L + n_R) \end{array}$$

Spin-1/2 XY chain (compactified boson CFT analysis)

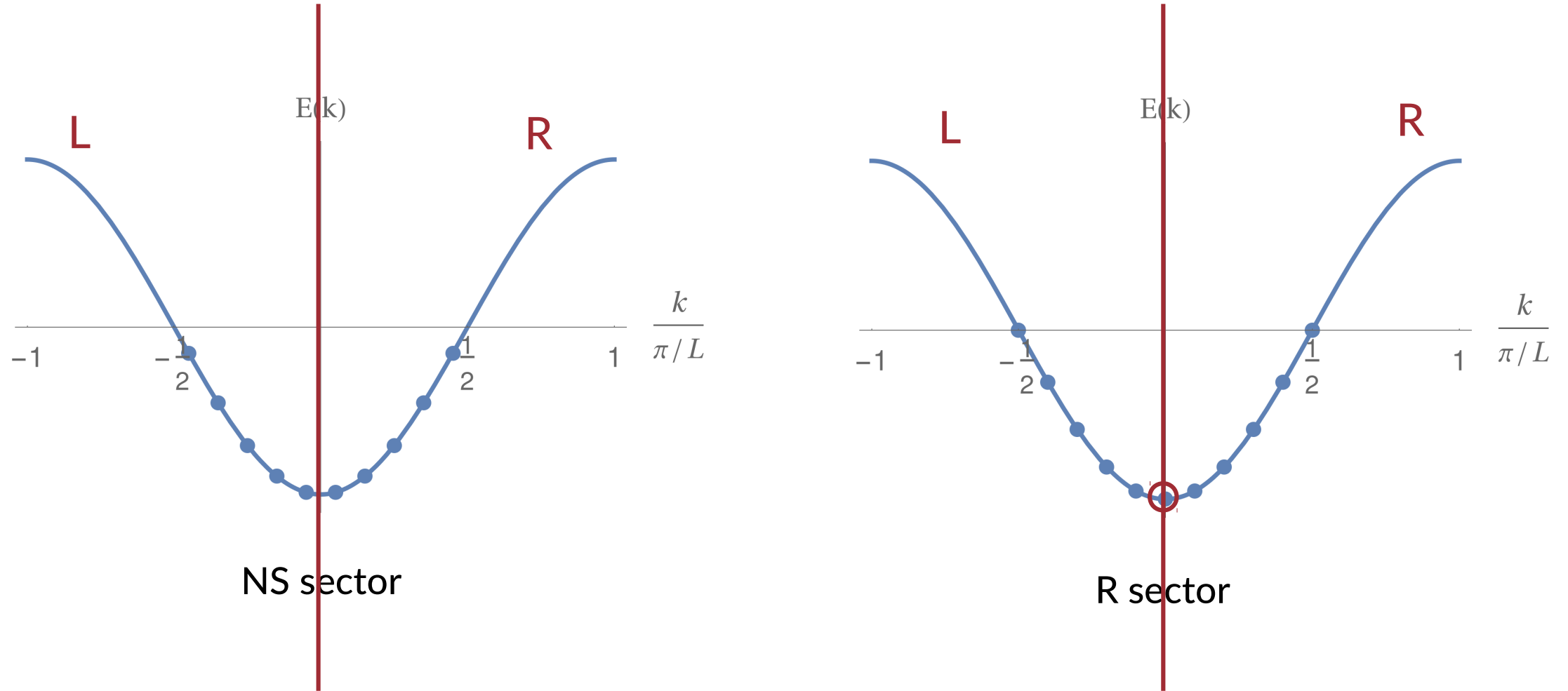


NS sector



R sector

Spin-1/2 XY chain (compactified boson CFT analysis)



Spin-1/2 XY chain (compactified boson CFT analysis)

(spinless) Luttinger liquid

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Responsible for the discontinuity at the boundary

Dual fields

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$$R = 2$$

$$R = 1$$

Spin-1/2 XY chain (compactified boson CFT analysis)

(spinless) Luttinger liquid

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$$\begin{aligned} R &= 2 \quad \text{😊} \quad \rightarrow \quad S_{\text{KB}} = \ln R = \ln 2 \\ R &= 1 \end{aligned}$$

Spin-1/2 XXZ chain

$$H = - \sum_{i=1}^L (S_i^x S_{i+1}^x + S_i^y S_{i+1}^y) + \Delta \sum_{i=1}^L S_i^z S_{i+1}^z \longrightarrow \Delta \sum_i (f_i^\dagger f_i - \frac{1}{2})(f_{i+1}^\dagger f_{i+1} - \frac{1}{2})$$

Density-density
interaction

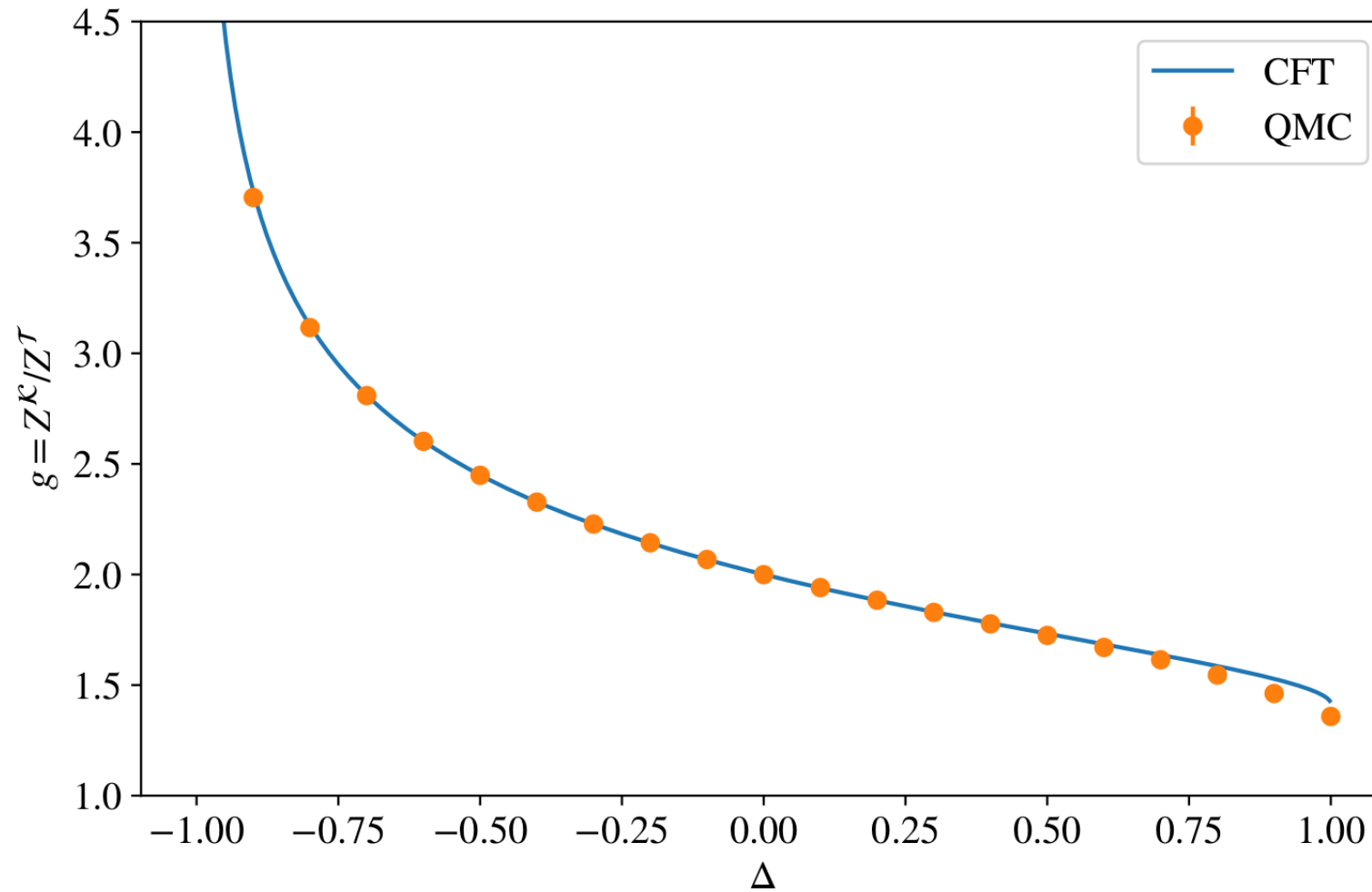
Umklapp
process

irrelevant for
 $-1 < \Delta < 1$

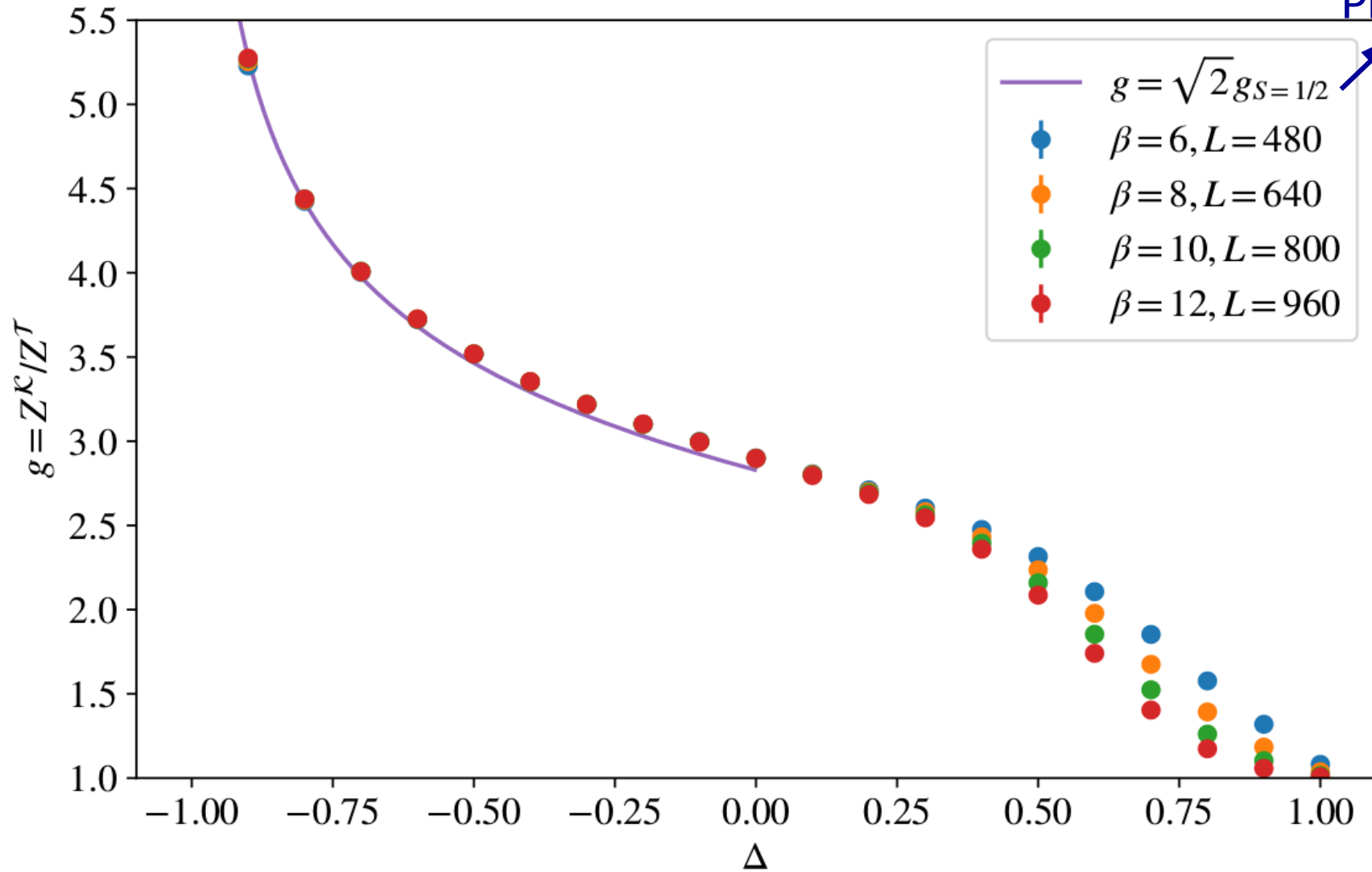
$$\begin{aligned} \phi &\Rightarrow \Phi = \sqrt{K} \phi \\ \theta &\Rightarrow \Theta = \theta / \sqrt{K} \end{aligned} \quad \longrightarrow \quad S_{\text{KB}} = \ln R = \ln(2\sqrt{K})$$

K Luttinger parameter

Spin-1/2 XXZ chain



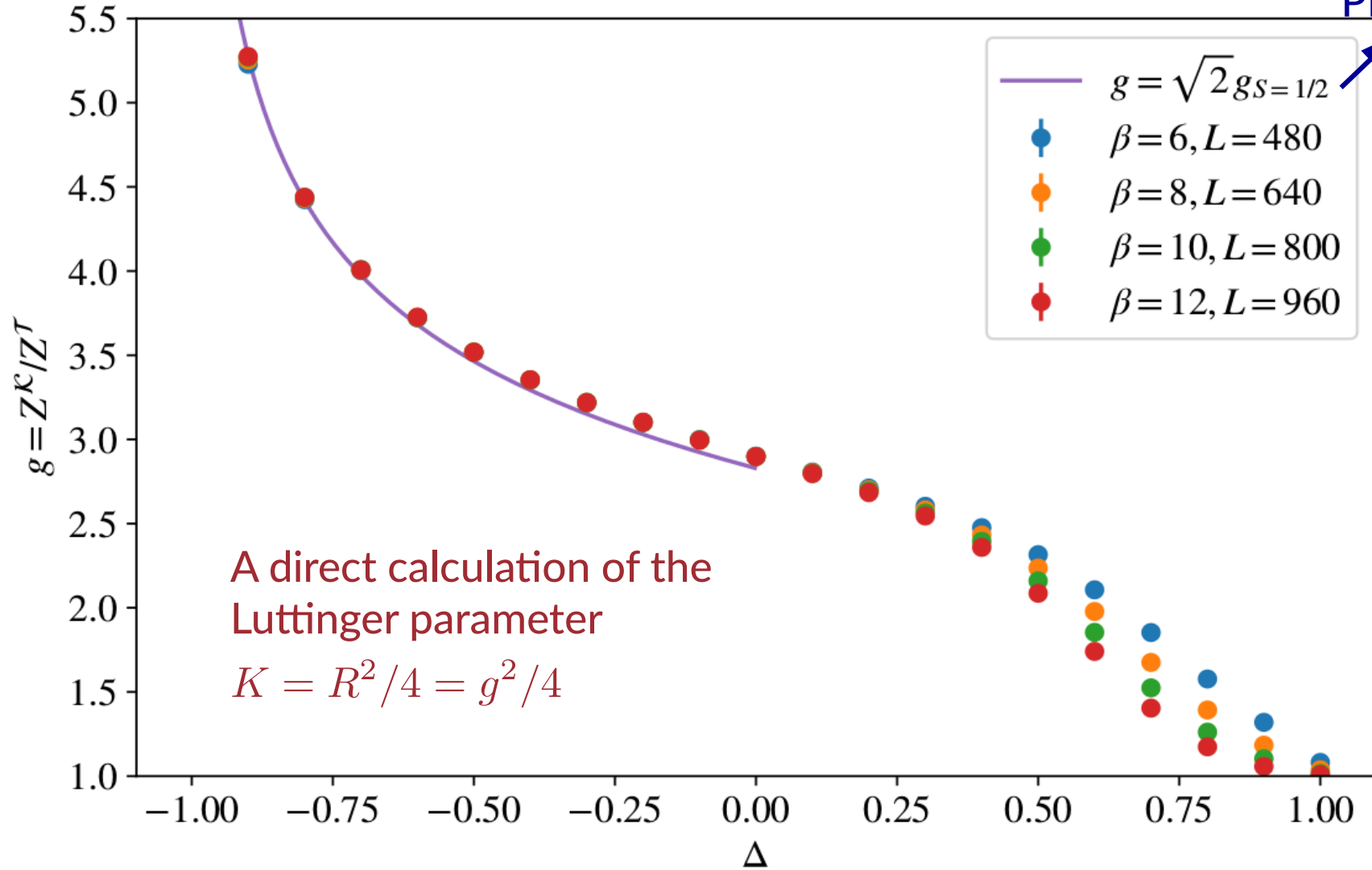
Spin-1 XXZ chain



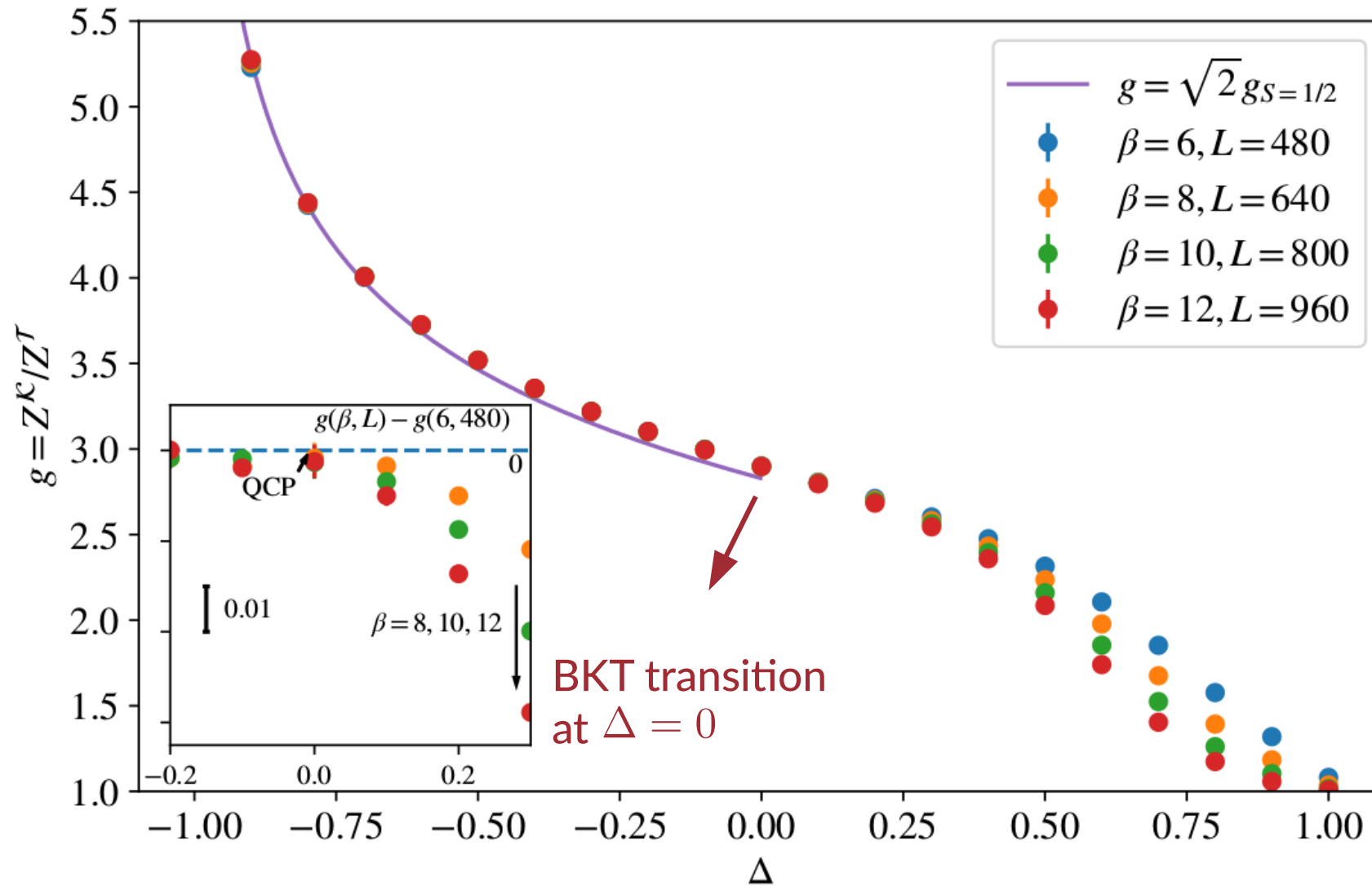
Alcaraz and Moreo,
PRB 1992

Spin-1 XXZ chain

Alcaraz and Moreo,
PRB 1992



Spin-1 XXZ chain



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Summary

- The Klein bottle entropy can be a powerful tool to characterize the CFT in conformal critical systems, especially in the rational CFT and the compactified boson CFT, where the Klein bottle entropy can be explicitly calculated
- In lattice models, with a careful definition and analysis of the lattice operation, the Klein bottle entropy can be efficiently extracted by our QMC method

Thank you for your attention!