

Several points:

1714年，华伦温标。

热学走向实验科学

热质说：希腊火元素学说的进一步发展

但是，不能解释摩擦生热。

相对立的学说：热是一种运动的表现形式

Francis Bacon (1561 - 1626)

强调理论必须根据实验事实，

热是一种运动

JOHNSON, 1711 - 1765

热是分子运动的表现，运动相互

Count Rumford, 1753 - 1814

制造花炮切下碎屑温度高，高温碎屑不断产生，
并非是一种运动不可。 (1798)

Humphry Davy, 1778-1829, chemist

而液体相互摩擦，产生融化

热功当量：

德国医生 Julius Robert Mayer, 1814-1878

能量守恒，也是能的一种形式，可与
机械能相互转化。

直接实验证据：

James Prescott Joule, 1818-1889

从 1840 年起，用电的热效应

1842 年起，机械生热法，

1850 年，得到科学界公认，能量守恒

— 热力学第一定律

准静态过程：

在过程进行中的每一步，物体都处在平衡态

— 随意、无摩擦、无限慢、可逆过程

可逆过程：相反方向进行，不对外界引起变化

资本主义，“永动机”

1775，巴黎科学院宣布了不再受理这类
动机的发明。

绝热过程：状态改变完全是由于机械能或

能的直接作用结果。（物理）

adiabatic : during the change of state,
no addition or removal of heat takes place;
that is, the system is isolated by
adiabatic walls (walls which do not
conduct heat) ($\delta Q = 0$)

热质沉 \rightarrow 热生 $\rightarrow C$

$C(T)$

标准温度 15°F

一克纯水在一个大气压下 $19.5^{\circ}\text{C} \rightarrow 15.5^{\circ}\text{C}$

所布热量

internal energy: 内动能 + 内势能

$$\Delta E_{in} = W_{ex}$$

$$\Delta U = W_{micro} + W_{macro}$$

$$= Q - W \text{ (here, } \exists \text{ by heat)}$$

Q : mysterious energy

$$\Delta U = Q - W$$

$$Q = \Delta U + W$$

计算公式: W , ΔU , Q .

Q 直接给

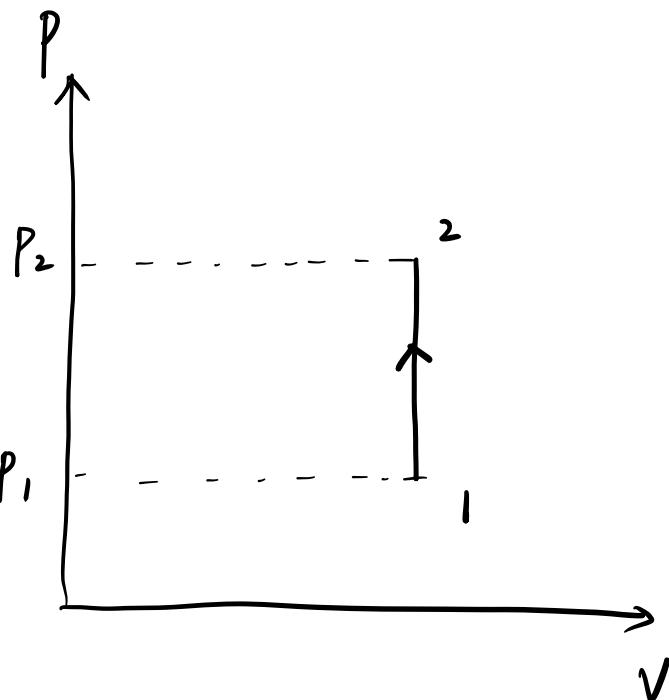
\cancel{W} $\Delta U + W$

Calculation :

I° 等容

isochore

1' air isochore P_1



$$\Delta W = \int p dV = 0$$

$$\Delta U = \frac{1}{2} n R (T_2 - T_1)$$

$$= \frac{1}{2} (P_2 V_2 - P_1 V_1) = \frac{1}{2} (P_2 - P_1) V$$

$$= C_V n \Delta T$$

$$\boxed{\frac{1}{n} \frac{\Delta U}{\Delta T} = C_V}$$

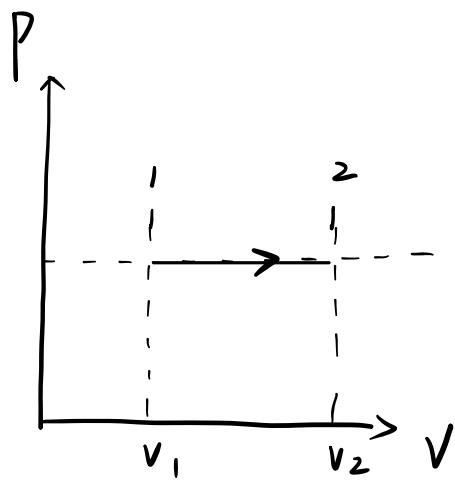
摩尔等容热容

$$\Delta Q = \Delta W + \Delta U = C_V n \Delta T$$

II. 等压:

isobaric

1, aison' barenic)



$$v_2 = 2v_1$$

$$\Delta w = F \Delta l = p s \Delta l = p \Delta V$$

$$\Delta w = \int_1^2 p dv = p (v_2 - v_1)$$

对于理想气体，内能就是内动能

$$U = \frac{1}{2} n R T$$

$$\Delta U = \frac{1}{2} n R (T_2 - T_1) = \frac{1}{2} n R (T_2 - T_1)$$

$$\downarrow \quad p_1 v_1 = n R T_1$$

$$= \frac{1}{2} p (v_2 - v_1)$$

焦耳热压热容

$$\Delta Q = \Delta U + \Delta w$$

$$= \frac{i+2}{2} p (v_2 - v_1)$$

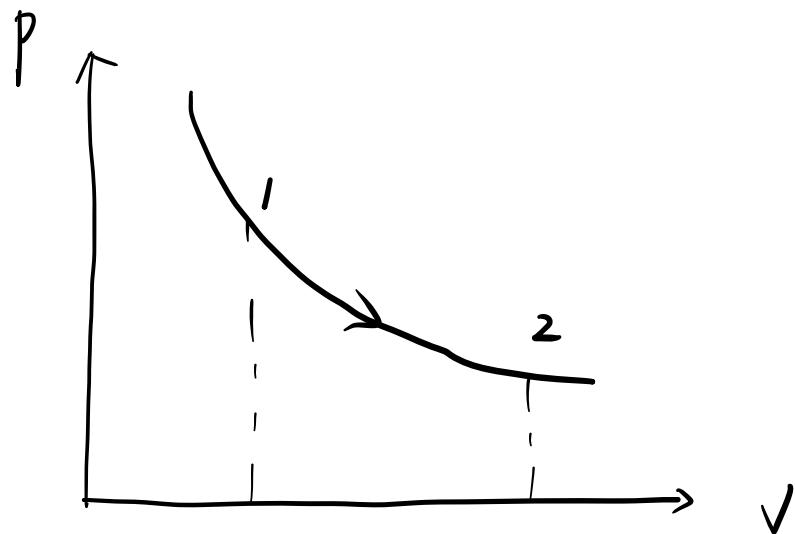
$$= \frac{i+2}{2} n R \Delta T$$

$$\frac{\Delta Q}{\Delta T} = \boxed{C_p} n R$$

$$\boxed{C_V + 1}$$

$$\begin{aligned}
 \text{Shankar's trick : } \Delta Q &= \Delta U + \Delta W \\
 &= \Delta U + P \Delta V \\
 &= \Delta U + \frac{\Delta (PV)}{nR\Delta T}
 \end{aligned}$$

III ^o Isothermal process



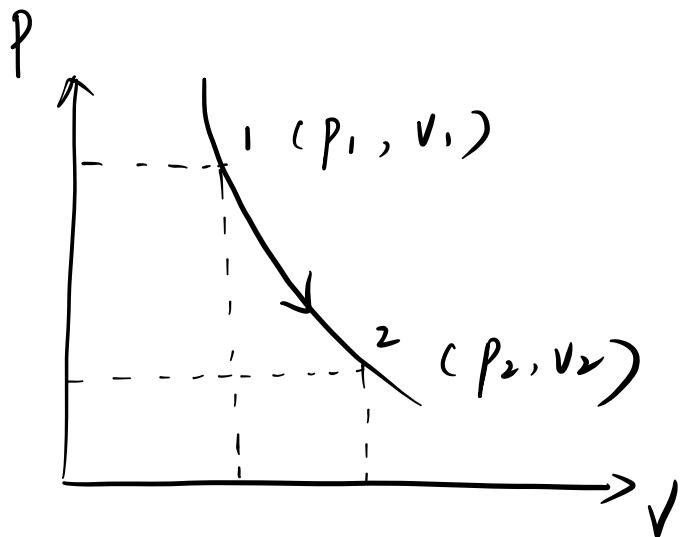
$$\begin{aligned}
 \Delta W &= \int_1^2 P \, dV \\
 &= \int_1^2 \frac{nRT}{V} \, dV \\
 &= nRT \ln \left(\frac{V_2}{V_1} \right)
 \end{aligned}$$

$$\Delta U = 0$$

$$\Delta Q = \Delta U + \Delta W = nRT \ln \left(\frac{V_2}{V_1} \right)$$

IV adiabatic process

$$\Delta Q = \Delta U + \Delta W = 0 \quad (\text{绝热})$$



(绝热材料 +
抽真空
(失去热交换))

$$\begin{aligned}\Delta U &= \frac{1}{2} n R (T_2 - T_1) = \int_{T_1}^{T_2} C_V dT \\ &= \frac{1}{2} (P_2 V_2 - P_1 V_1)\end{aligned}$$

$$\frac{1}{2} n R dT + p dV = 0$$

$$\frac{1}{2} (p dV + V dp) + p dV = 0$$

$$\frac{i+2}{2} p dV \quad \gamma p dV = -V dp$$

$$= -\frac{i}{2} V dp$$

$$\gamma \frac{dV}{V} = -\frac{dp}{p}$$

(*)

$$\gamma = \frac{i+2}{i} = \frac{C_p}{C_v}$$

$$\gamma d\ln V = -dmp + C$$

$$\gamma \int_{(P_1, V_1)}^{(P_2, V_2)} d\ln V = - \int_{(P_1, V_1)}^{(P_2, V_2)}$$

$$dmp + C_1$$

$$\gamma \ln \frac{V_2}{V_1} = - \ln \frac{P_2}{P_1} + C_1$$

$$\left(\frac{V_2}{V_1} \right)^\gamma = C_2 \frac{P_1}{P_2}$$

$$P_1 V_1^\gamma = P_2 V_2^\gamma = \text{const}$$

directly start from $(*)$

$$\gamma p dv + v dp = 0$$

$$\gamma \frac{dv}{v} + \frac{dp}{p} = 0 \Rightarrow \gamma d\ln v + d\ln p = 0$$

$$\gamma \ln V + \ln p = C \rightarrow PV^\gamma = \text{const}$$

$$\Delta U = \frac{1}{2} nR (T_2 - T_1)$$

$$= n C_V (T_2 - T_1)$$

$$\therefore \frac{nR}{\gamma-1} (T_2 - T_1)$$



$$\gamma = \frac{C_P}{C_V} = \frac{\frac{1}{2} + 1}{\frac{1}{2}}$$

$$\frac{\gamma-1}{1} = \frac{C_P - C_V}{C_V} = \frac{\frac{1}{2} + 1 - 1}{\frac{1}{2}}$$

$$= \frac{R}{C_V} \Rightarrow$$

$$C_V = \frac{R}{\gamma-1}$$

$$dU = \int p dV$$

$$= \int \frac{C}{V^\gamma} dV$$

$$= C \frac{\gamma^{-\gamma+1}}{-\gamma+1} \Big|_1^2$$

$$= C \frac{V_2^{1-\gamma} - V_1^{1-\gamma}}{1-\gamma}$$

$$\downarrow \quad PV^\gamma = C$$

$$PV = nRT$$

$$V^{\gamma-1} = \frac{C}{nRT}$$

$$= C \frac{\left[\frac{nRT_2}{C} - \frac{nRT_1}{C} \right]}{1-\gamma}$$

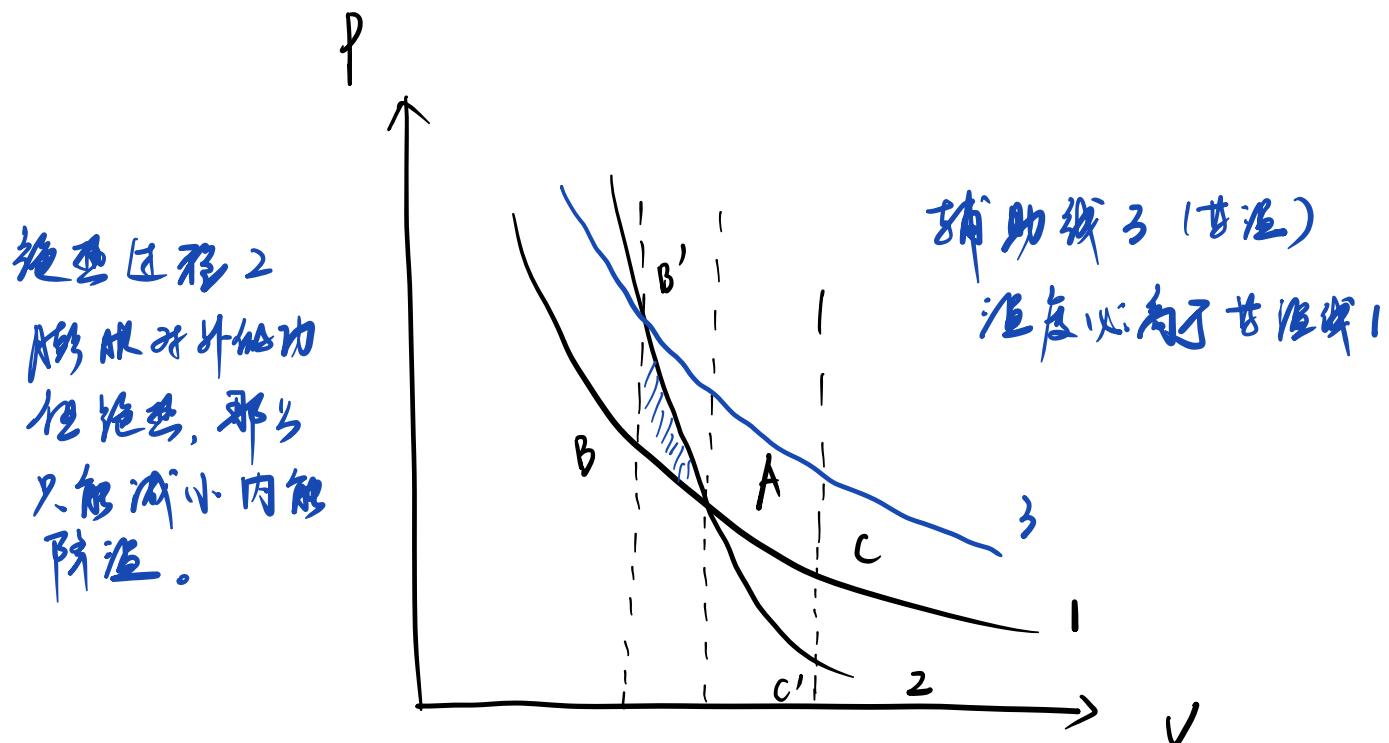
$$= \frac{nR}{1-\gamma} (T_2 - T_1) = -\Delta U$$

$$\Delta H = \Delta U + \Delta W = 0 \quad (\text{绝热})$$

Isothermal

vs

adiabatic



Which one is isothermal ?

$$\Delta F = \Delta U + \Delta W$$

$$\left\{ \begin{array}{l} B' \rightarrow A \\ B \rightarrow A \end{array} \right. \quad T_{B'} > T_B$$

(pV = nRT)

绝热： $\Delta U + \Delta W = 0$ 更多的作功妨碍

内能系数

$$\begin{cases} A \rightarrow C \\ A \rightarrow C' \end{cases}$$

对子绝热： $d\varphi = du + pdv = 0$

膨胀对外做功
即：降内能， $T \downarrow$

$$U_{C'} < U_C$$

$$T_{C'} < T_C$$

$$PV = c_1 \Rightarrow P = c_1 / V \quad \textcircled{1}$$

$$PV^\gamma = c_2 \quad P = c_2 / V^\gamma \quad \textcircled{2}$$

$$\textcircled{1} \quad \frac{\partial P}{\partial V} = c_1 (-1) V^{-2}$$

$$\textcircled{2} \quad \frac{\partial P}{\partial V} = c_2 (-\gamma) V^{(-\gamma-1)}$$

$$\textcircled{1} / \textcircled{2} = \frac{c_1}{c_2} \cdot \frac{1}{\gamma} (V^{\gamma-1} = \frac{c_2}{c_1})$$

$$= -\frac{1}{\gamma}$$

因为负数

$$\left| \frac{\partial P}{\partial V} \right|_{\textcircled{1}} < \left| \frac{\partial P}{\partial V} \right|_{\textcircled{2}}$$

绝热及限峭

heat capacity

$$\left\{ \begin{array}{l} C_V = \frac{1}{n} \left(\frac{dU}{dT} \right)_V \quad \text{摩尔等容热容} \\ C_P = \frac{1}{n} \left(\frac{dH}{dT} \right)_P \quad \text{摩尔等压热容} \end{array} \right.$$

$$H = U + PV \quad (\text{enthalpy} = \text{internal energy} + PV)$$

$$dQ = dU + dW$$

$$dQ = dU + PdV$$

(the differential form of 1st law)

$$C_V = \frac{1}{n} \left(\frac{dU}{dT} \right)_V$$

ideal gas

$$U = \frac{3}{2} N k_B T = \frac{3}{2} n R T$$

$$\left(\frac{dU}{dT}\right)_V = \frac{3}{2} nR$$

$$C_V = \frac{3}{2} R$$

$$PV = nRT$$

$$V = \frac{nRT}{P}$$

$$\frac{dV}{dT} = \frac{nR}{P}$$

$$dQ = dU + PdV \quad C_p \equiv \frac{1}{n} \left(\frac{dQ}{dT} \right)_P$$

$$C_p = \frac{1}{n} \left(\frac{dU}{dT} \right)_P + \frac{1}{n} P \left(\frac{dV}{dT} \right)_P$$

$$C_p = \left(\frac{dQ}{dT} \right)_P = C_V + \frac{P}{n} \times \frac{nR}{P}$$

$$dQ = dU + PdV$$

$$+ V \frac{dp}{T}$$

$$= C_V + R = \frac{5}{2} R$$

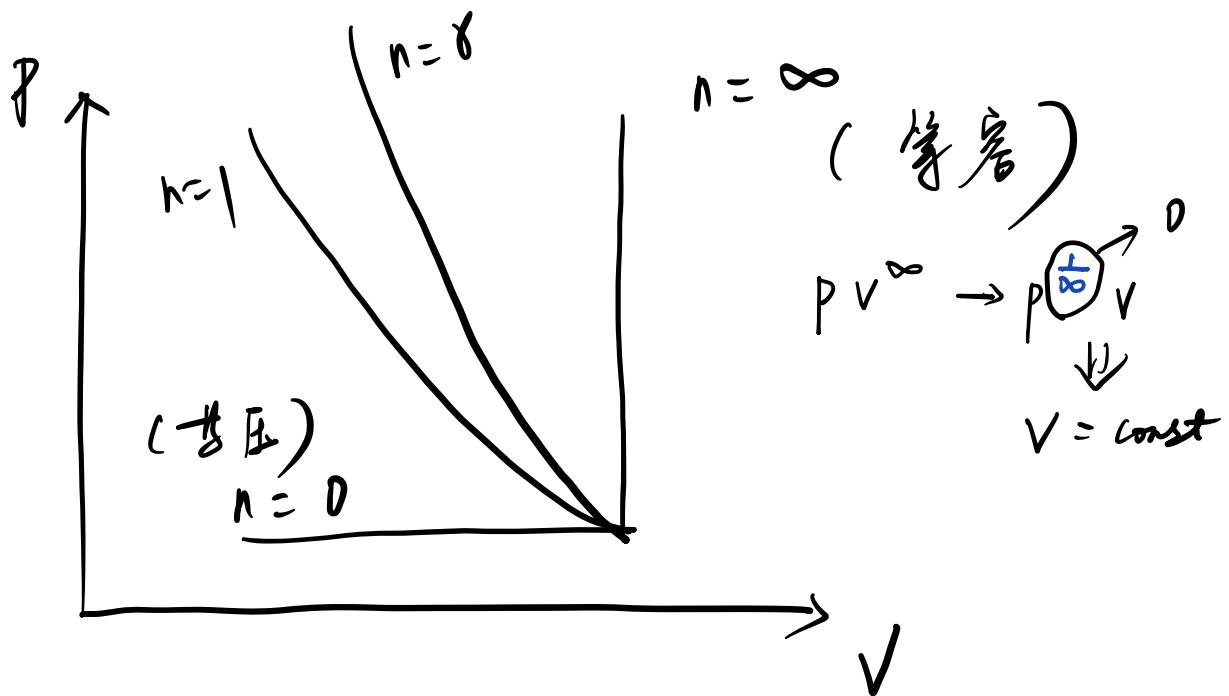
$$= dU + d(C_p V)$$

$C_p = C_V + R$

C_p vs C_V : experiment vs theory

多方过程 (polytropic process)

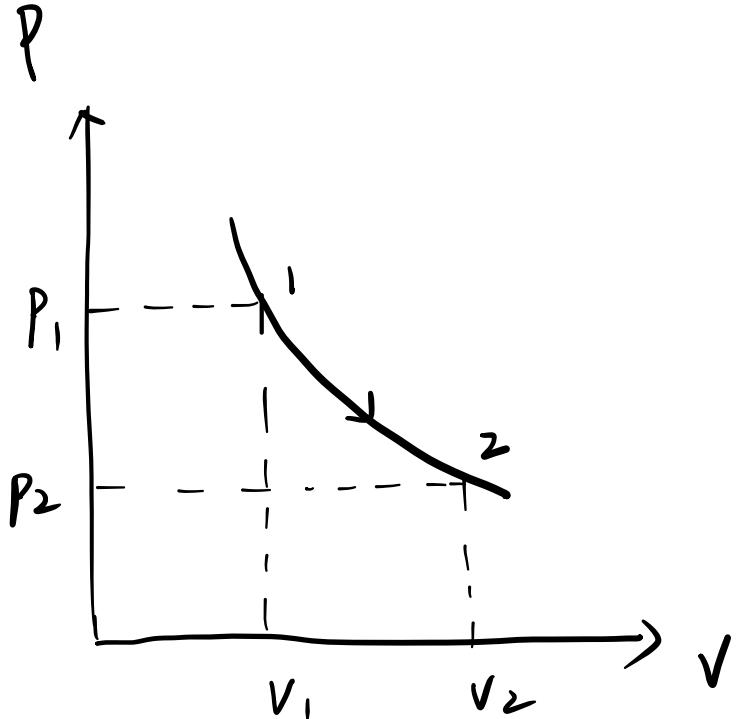
$$P V^n = \text{const}$$



$$n \in (1, \gamma)$$

实际过程

$$d\omega = P dV$$



$$dQ = dU + dW$$

$$\Delta W = \int_1^2 P dV$$

$$= \int_1^2 \frac{c}{V^n} dV = \frac{V^{1-n}}{1-n} \Big|_{V_1}^{V_2}$$

$$= c \frac{V_2^{1-n} - V_1^{1-n}}{1-n}$$

$$dU = \frac{1}{2} n R dT = \frac{1}{2} d(PV)$$

$$\Delta U = \frac{1}{2} (P_2 V_2 - P_1 V_1)$$

$$= \frac{i}{2} \left(-\frac{c}{V_2^{n-1}} - \frac{c}{V_1^{n-1}} \right)$$

$$= \frac{i}{2} c (V_2^{1-n} - V_1^{1-n})$$

$$\Delta W = \frac{1}{1-n} c (V_2^{1-n} - V_1^{1-n})$$

$$\Delta U = \frac{i}{2} c (V_2^{1-n} - V_1^{1-n})$$

$$\Delta Q = \left(\frac{i}{2} + \frac{1}{1-n} \right) c (V_2^{1-n} - V_1^{1-n})$$

$$\xrightarrow{\gamma = r} \frac{1}{1-r} = \frac{1}{1 - \frac{c_p}{c_v}} = \frac{c_v}{c_v - c_p}$$

$$\frac{c_v}{c_v - c_p} = \frac{i}{-2}$$

↑

$$\frac{c_v}{c_p} = \frac{i}{i+2}$$

多方 → 绝热

$$\frac{C_p}{C_V} = \frac{5}{3}$$

(单原子分子)

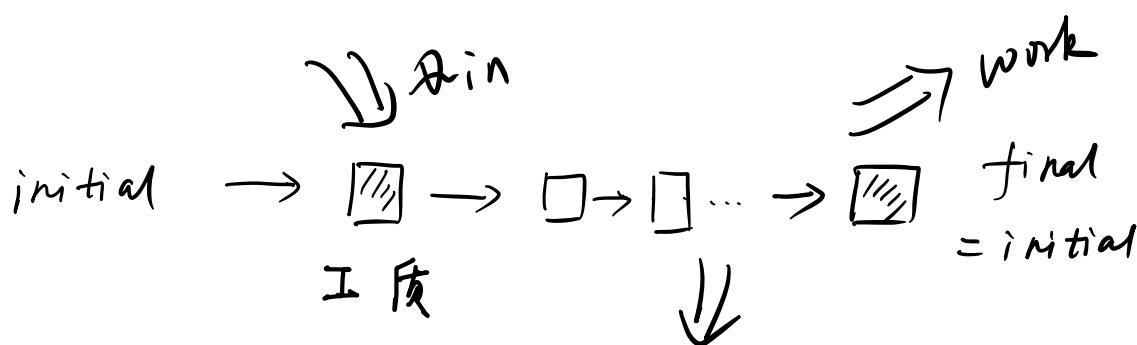
(C_p : 体积膨胀，做功，吸
热量)

otto cycle in thermal physics II

Lin Hsin-Han (现在网站课程更名为
热统计物理)

heat engine : A device

converts "heat" into "work"



Q_{out} (waste heat)